



# **Prospects for He Pumping by Nano-Bubble Formation in Liquid Metals**

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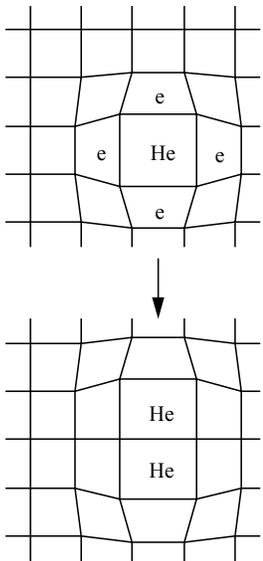
Outline:

- Sandia's Nano-Bubble Evolution code for He in metals
- Modifications for liquid metals
- Sample calculations for He implantation in liquid Ga, Sn, Li
- Implications for He pumping
- Experimental validation

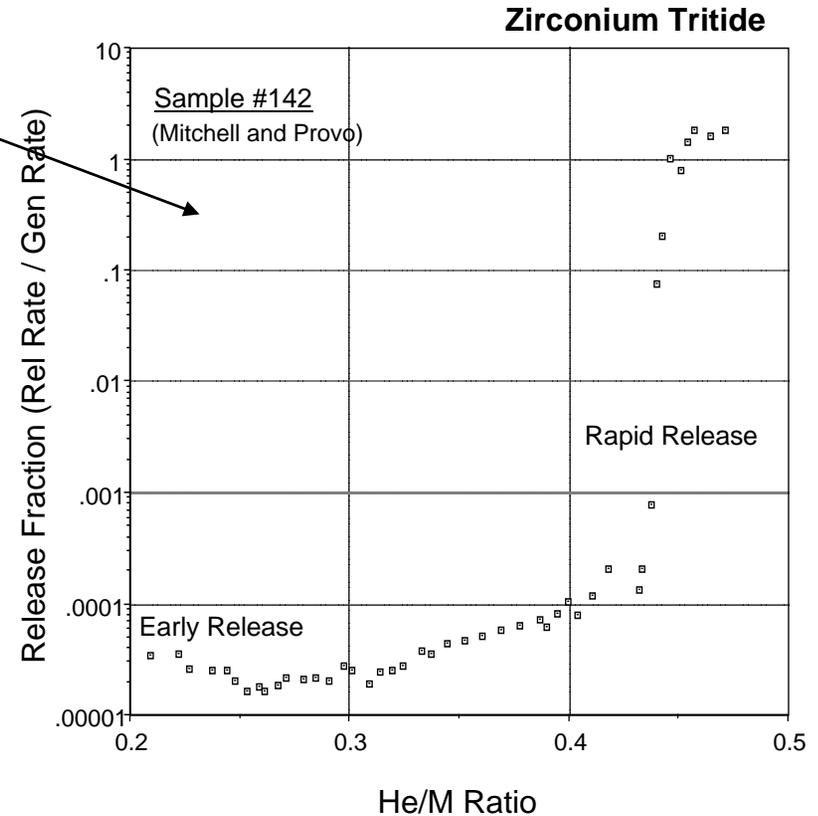
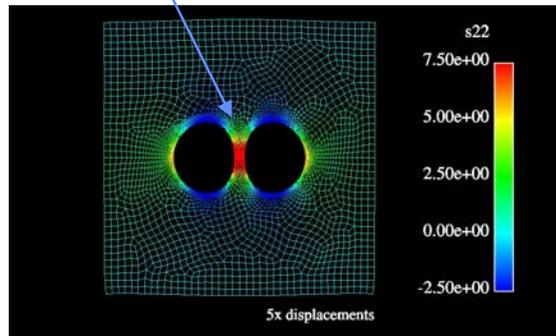


# Sandia's Nano-Bubble Evolution (NBE) code models the retention of He generated in metals.

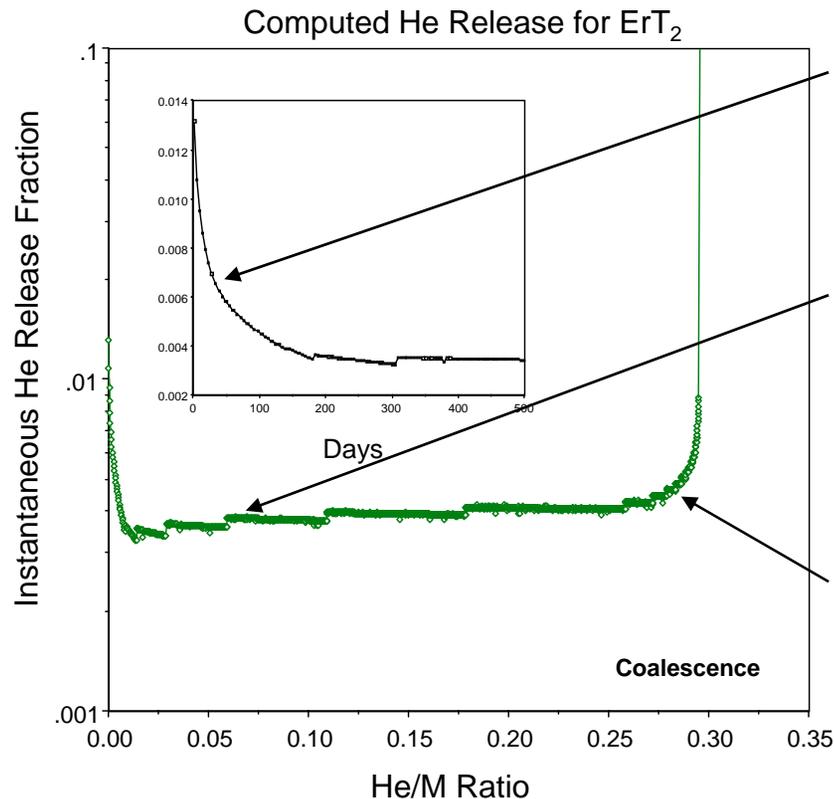
- The NBE code was developed to model  $^3\text{He}$  retention in aging solid metal tritides.



- It follows He bubble nucleation, growth, and release by inter-bubble ligament fracture.



# The computed He release for solid tritides has all the characteristics of observed release.



- Initially, release is high until bubbles become large enough to compete with nearby surfaces or grain boundaries.
- Mobile He near surfaces & GB's produce the Early Release Fraction.
  - which slowly increases with the “breaching” of near-surface bubbles.
- Rapid Release of He occurs when the growing bubble network becomes sufficiently interconnected.

*Model shows how material parameters affect each part of release spectrum.*

# For Liquid Metals, we modify the conditions for bubble stability and add bubble motion.

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- The nano-bubble stability is determined by the liquid metal's surface tension  $\gamma$ :

$$\text{Binding energy of N-th He, } E_N = 4\pi [ (r_{N-1}^2 + r_1^2) - r_N^2 ] \gamma.$$

- The bubble pressure  $p_N$  and radius  $r_N$  are related by

$$p_N = 2\gamma/r_N.$$

- We assume spherical bubbles so that

$$V_{\text{bubble}} = (4/3)\pi r_N^3 = n v_a,$$

where  $v_a(p,T)$  is the He atomic volume as determined by the bulk He equation-of-state, corrected for bubble wall curvature effects.

- Nano-bubble migration is included using the Stokes-Einstein equation for diffusion of sub-micron particles:

$$D_N = kT/6\pi\eta r_N.$$



# Concentration profiles of diffusing He atoms and n-atom bubbles are computed.

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- Coupled differential equations follow He *atom* and *bubble* concentrations as a function of depth and plasma exposure time:

$$dc_n/dt = - D_n d^2c_n/dx^2 + (\text{Source terms})_n - (\text{Loss terms})_n$$

where  $c_n$  = concentration of bubbles with N atoms.

Source terms: for  $n=1$ , implant flux =  $\phi(x,t)^{\text{TRIM}}$

n-species formation =  $c_1 c_{n-1}$ ,  $c_m c_{n-m}$ ,  $c_{n+1} \exp(-E_{n+1}/kT)$   
(growth, coalescence, dissociation)

Loss terms: n-species dissociation factors =  $c_n \exp(-E_n/kT)$

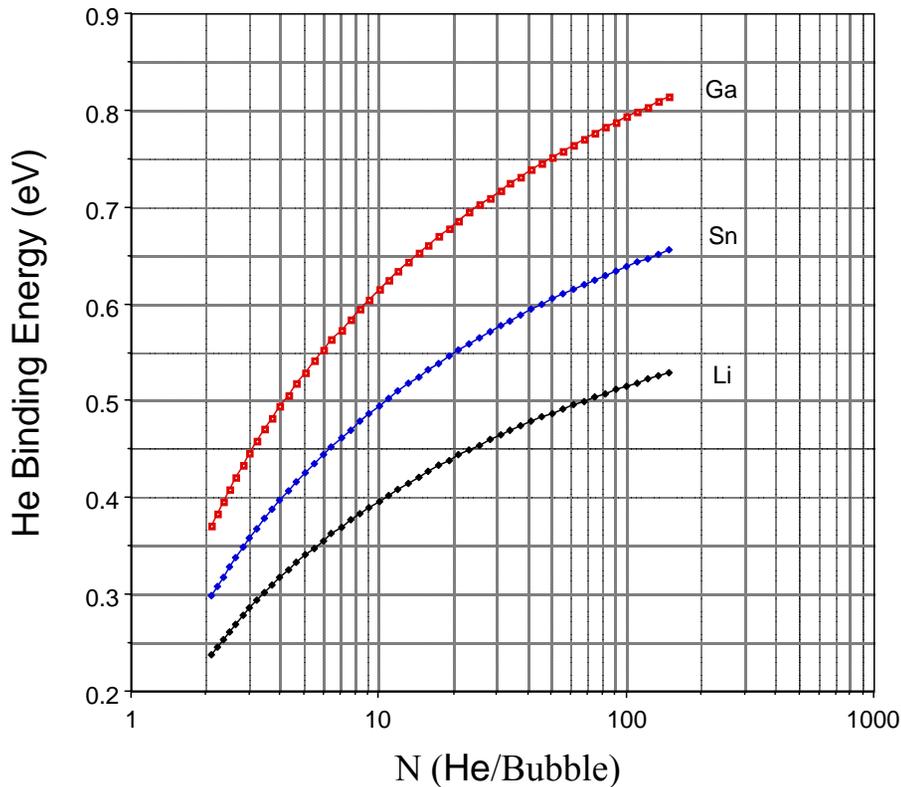
n-species promotion factors =  $c_1 c_n$ ,  $c_m c_n$

(Formation and loss rates also involve attempt frequency and bubble size factors.)

- Species bins increase geometrically:  $N = 1, 2, 4, 8, 16, \dots, 2^{n-1}$ .



# The model shows small bubbles rapidly dissociate due to a low He binding energy to the bubble.



- Generation of large stable bubbles from diffusing atoms requires:

species promotion > dissociation

$$s_n c_1 > N \exp(-E_N/kT),$$

$s_n$  = bubble surface area

$c_1$  = atomic He concentration.

- At .01 A/cm<sup>2</sup>, 600 K,  $c_1 \sim 10$  appm; which requires

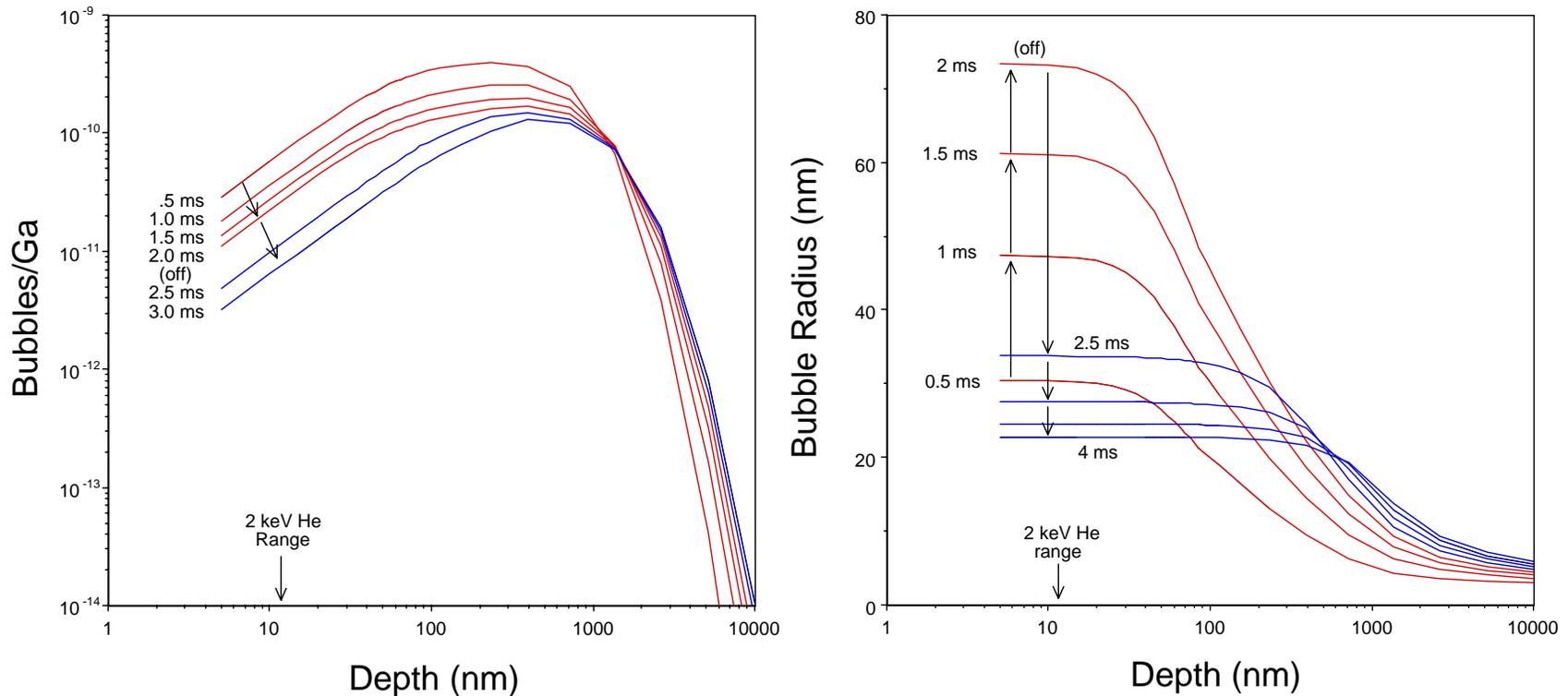
$$E_N \geq 0.4 \text{ eV.}$$

- Bubble promotion in Li requires a large  $c_1$  or very high He flux.

*The larger surface tensions for Sn and Ga rapidly produce bigger bubbles.*

# The profiles show bubbles are formed well beyond the He implant range.

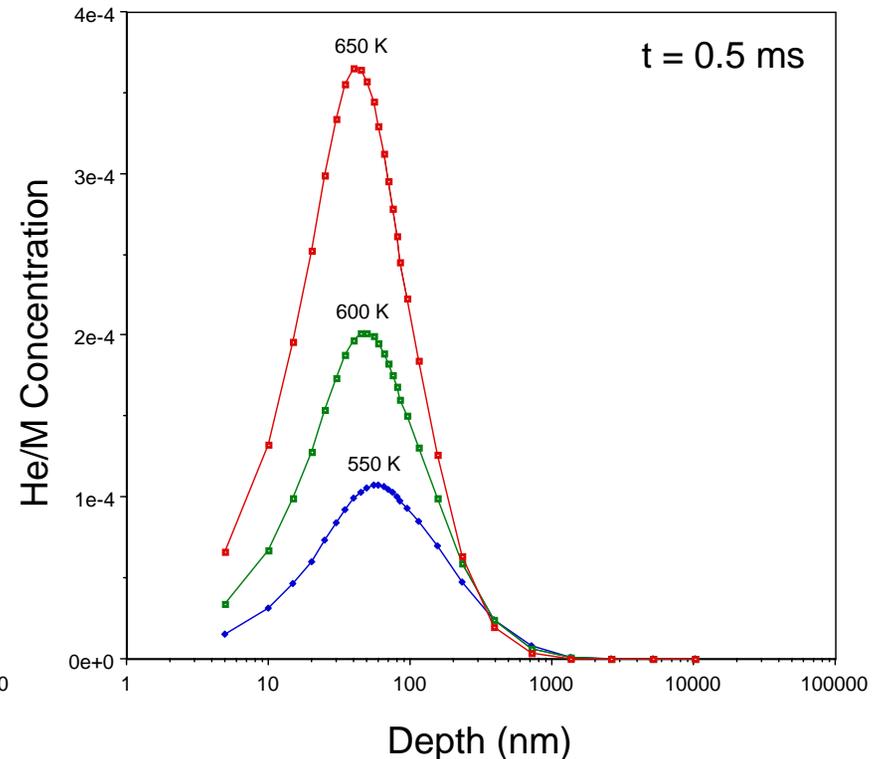
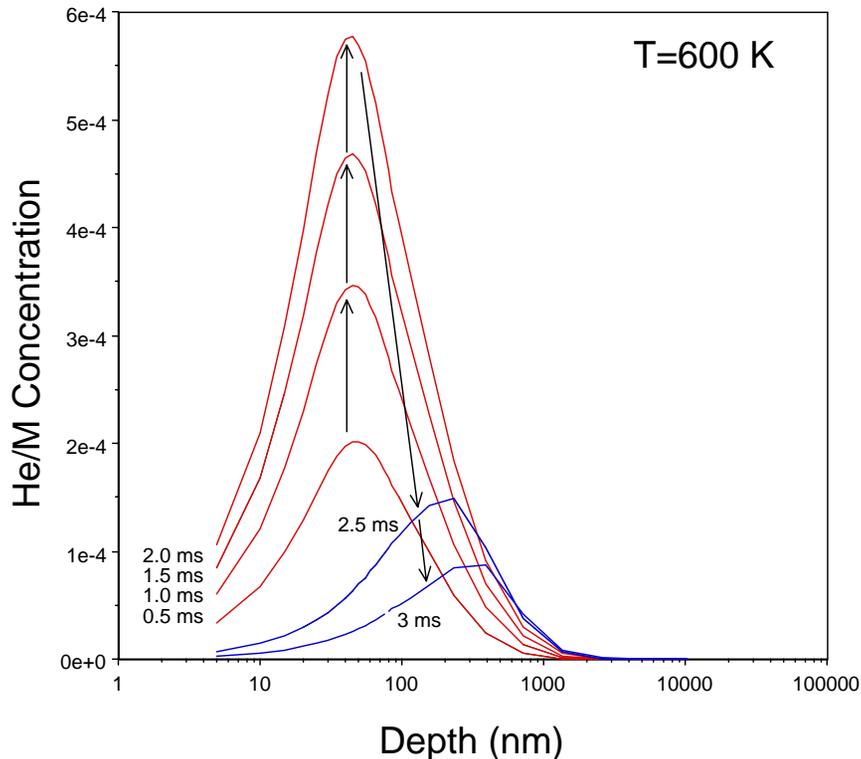
Liquid Ga: 600 K, 2 keV He,  $2e18$  He/cm<sup>2</sup>-s, 2 ms implant pulse



- Bubble concentrations drop with time as bubble sizes increase, due to bubble coalescence.

# Most of the He concentration remains within the surface region, and it increases with temperature.

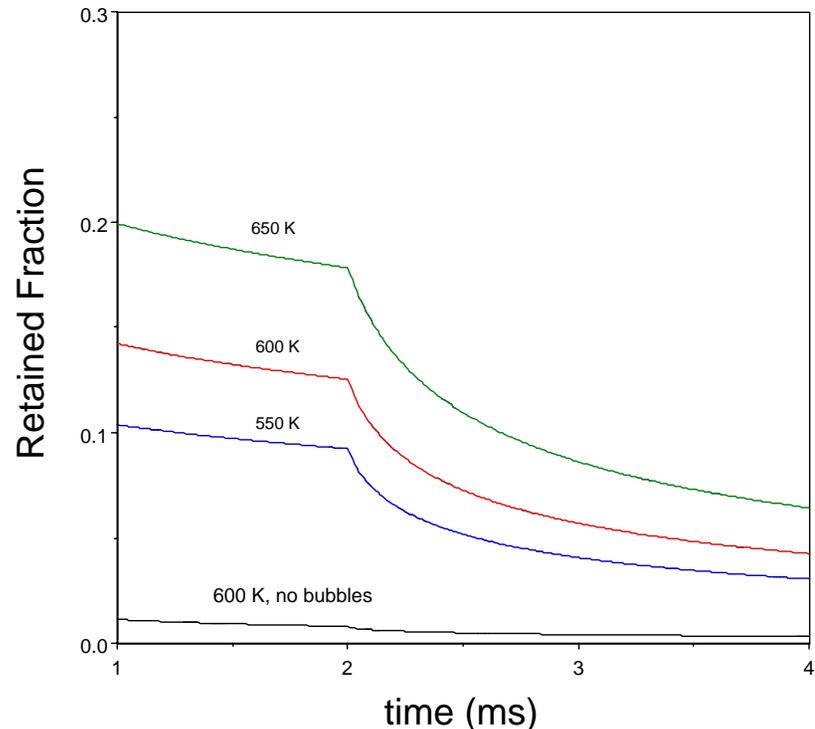
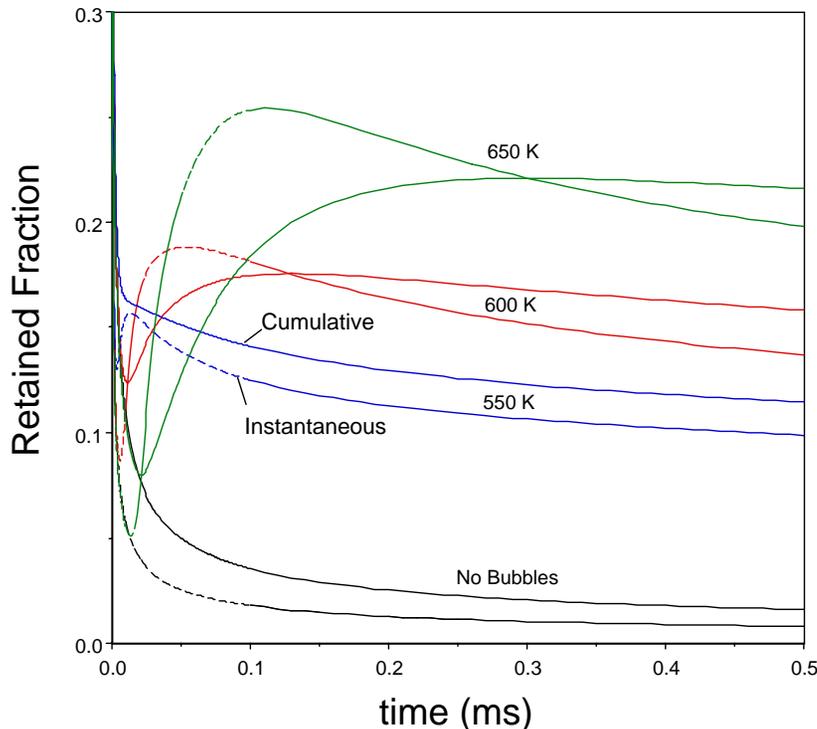
Liquid Ga: 2 keV He,  $2 \times 10^{18}$  He/cm<sup>2</sup>-s, 2 ms implant pulse



- Increasing the mobility promotes bubble growth and coalescence, creating larger bubbles.

# He retention in Ga reaches 10-20% over a 2 ms exposure pulse, then drops relatively slowly.

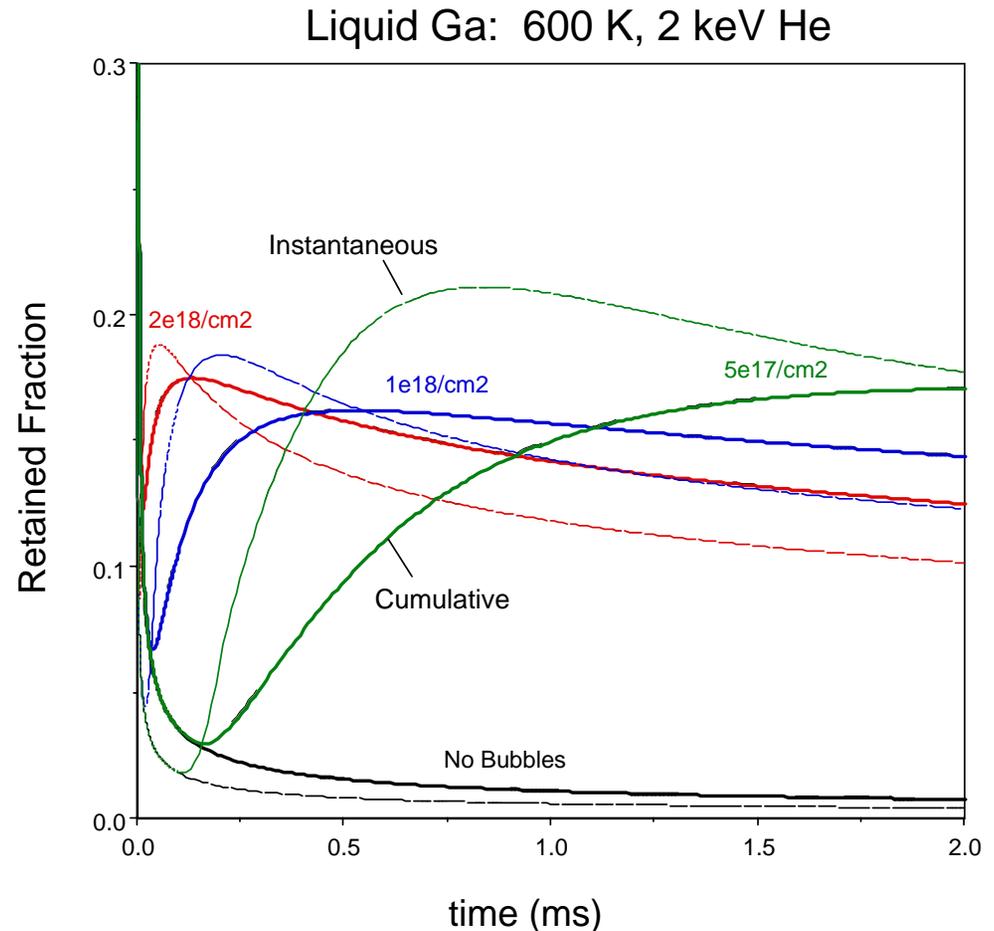
Liquid Ga: 600 K, 2 keV He,  $2 \times 10^{18}$  He/cm<sup>2</sup>-s



- The large quantity of He trapped in large bubbles in flowing Ga provides time for transport well into a capture duct.

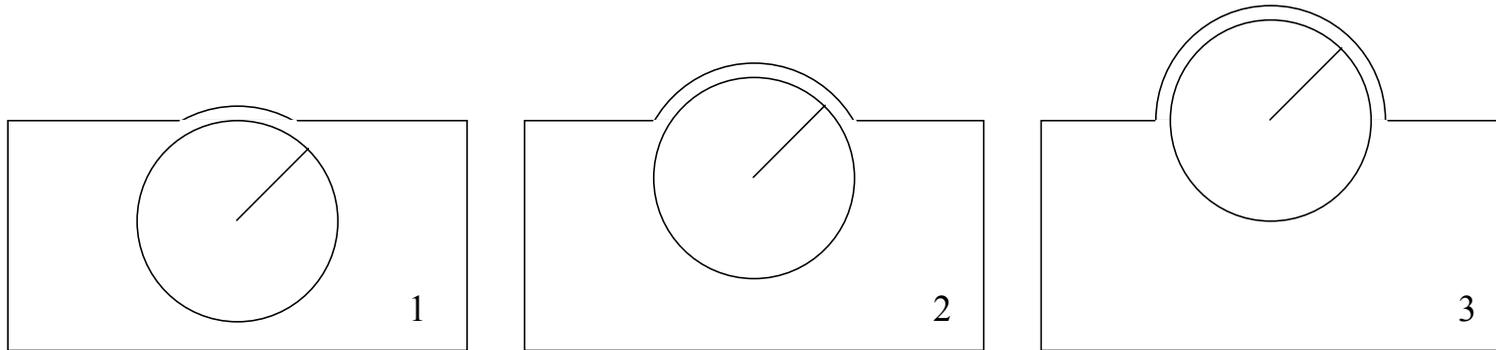
# The retention builds more slowly at lower He flux, but can reach similar levels.

- It takes longer to achieve a high concentration of large bubbles.
- At peak retention, there are increasingly fewer, but larger bubbles.
- For even lower fluxes, the retention drops as  $\mu\text{m}$ - $\text{mm}$  size bubbles breach by growth.
  - Collapse of the large cavities left behind can produce *undesirable* metal ejection -- *splattering*.



# He released from large bubbles depends on the condition for bubble “breach”.

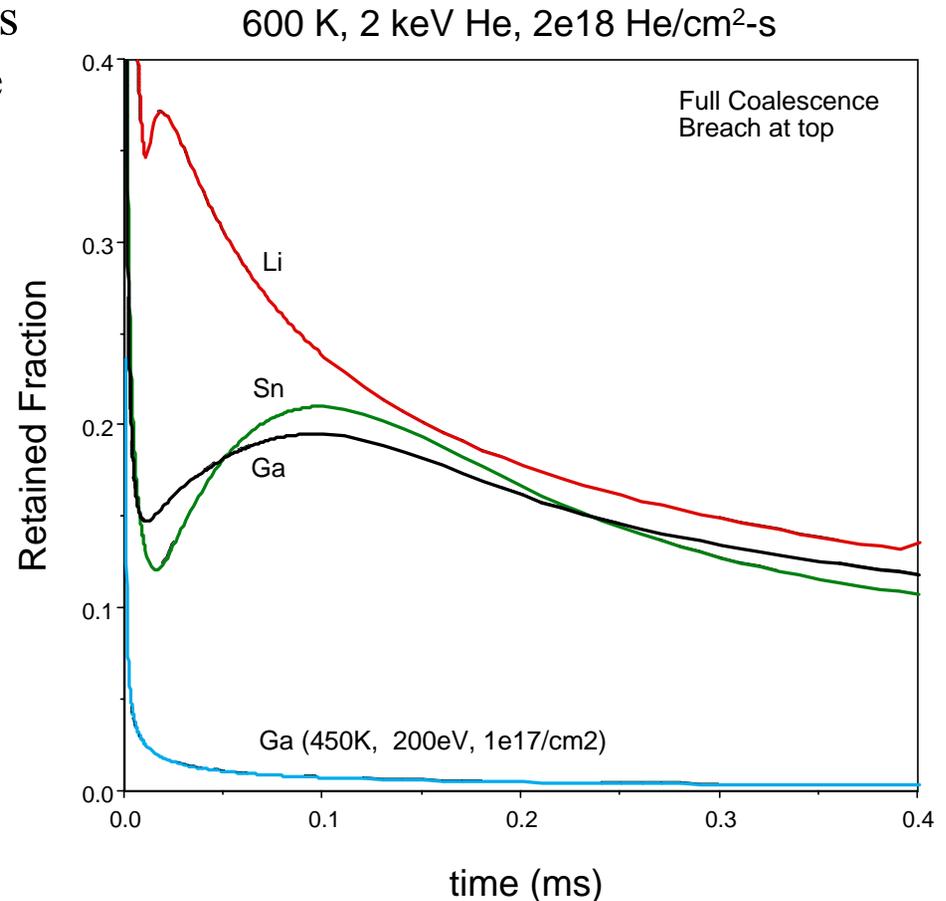
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- The previous calculations assume bubble breach at position 3.
- Breach may occur at some intermediate position 2.
  - Surface tension from the additional (outside) surface may shrink the bubble and increase its pressure.
  - Plasma heating of the bubble “skin” may accelerate the breach.
- Breach at position 1 should provide more conservative retention values.
  - Collapse of the resultant, deep cavity may produce metal ejection.
  - Is there a minimum bubble size for this to occur?

# With breach at the bubble top, retention is similar for Ga, Sn, and Li at .1-.2 ms

- For Li, very high concentrations of small bubbles ( $N \leq 256$ ) cause early, sustained high retention.
  - He is implanted deeper in Li.
  - The bubbles never get large.
  - Retention *depends strongly* on temperature and flux.
- For Ga and Sn, the retention grows until bubbles breach.
  - Breach of very large bubbles occurs by rapid coalescence -- not by diffusion.
  - Ret. is *less sensitive* to  $T, \phi$ .



# In summary, flowing Liquid Metals have potential for pumping He from plasma devices.

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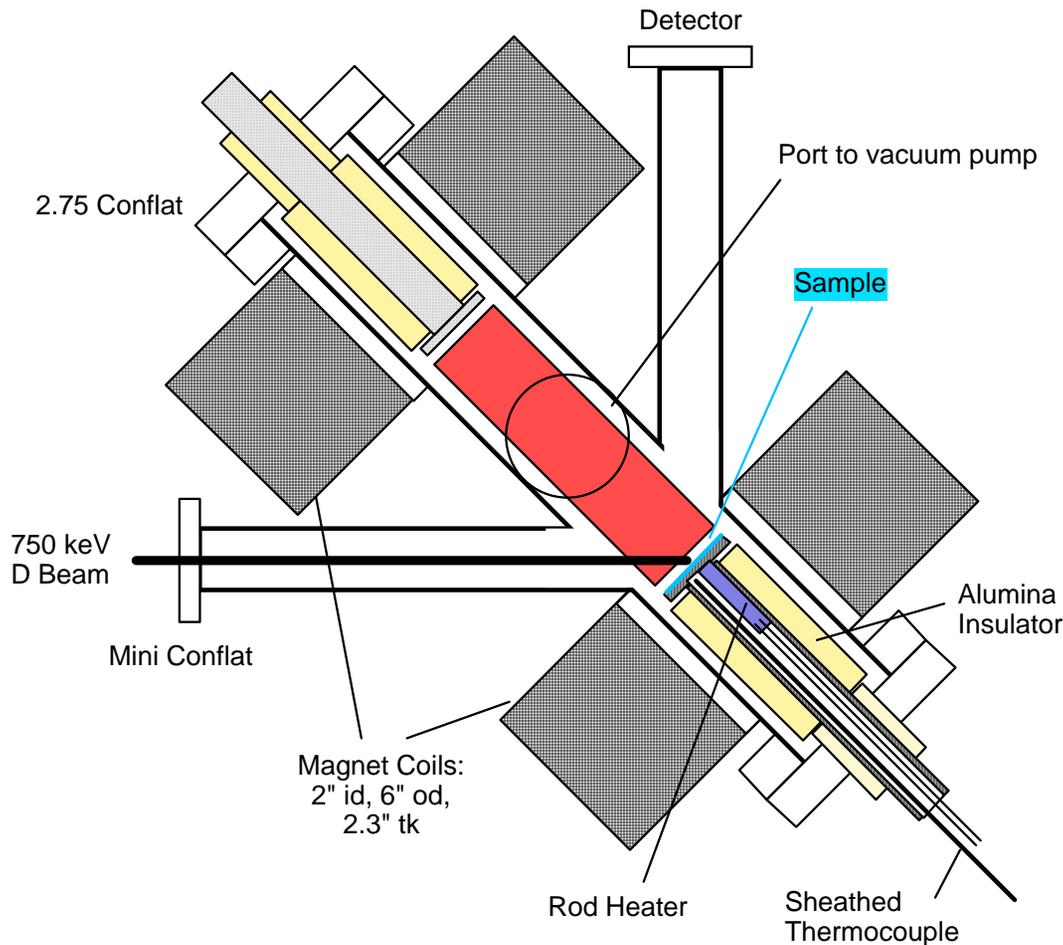
Based on very limited number of code runs -

- Retention is 10-20% over 1 ms pulse
  - For a given particle energy, the retention can be optimized using temperature, flux, and fluid flow.
- Best flux ranges:
  - Ga: low-moderate flux (.1-.2 A/cm<sup>2</sup>)
  - Sn: moderate-high flux (.2-.3 A/cm<sup>2</sup>)
  - Li: very high flux (>.3 A/cm<sup>2</sup>)
- Ga and Sn are less sensitive to variations in T, flux.
  - Net retention *increases* with temperature.
- Ga and Sn will provide better retention during transit into the collection duct.
  - But, Ga and Sn may splatter at high flux.

Future computations will examine the full parameter space and look at retention for mixed He-D plasmas.



# We hope to set up an experiment to test the model and provide a low flux benchmark of the code.



- Penning trap produces flux of  $10^{21}$  He/m<sup>2</sup>-s at 1 keV.
- Liquid Ga or Sn (1 cm<sup>2</sup>) covers trap (cathode) plate.
- <sup>3</sup>He added to plasma is *profiled in LiqM* by d<sup>+</sup> beam.
- High steady-state retention will signify bubble formation.
- Low retention (no bubbles) will provide a measurement of effective He diffusivity.