

Cross-field Plasma Transport in the Scrape-off Layer

**A key physics issue for SOL, divertor,
and core plasma**

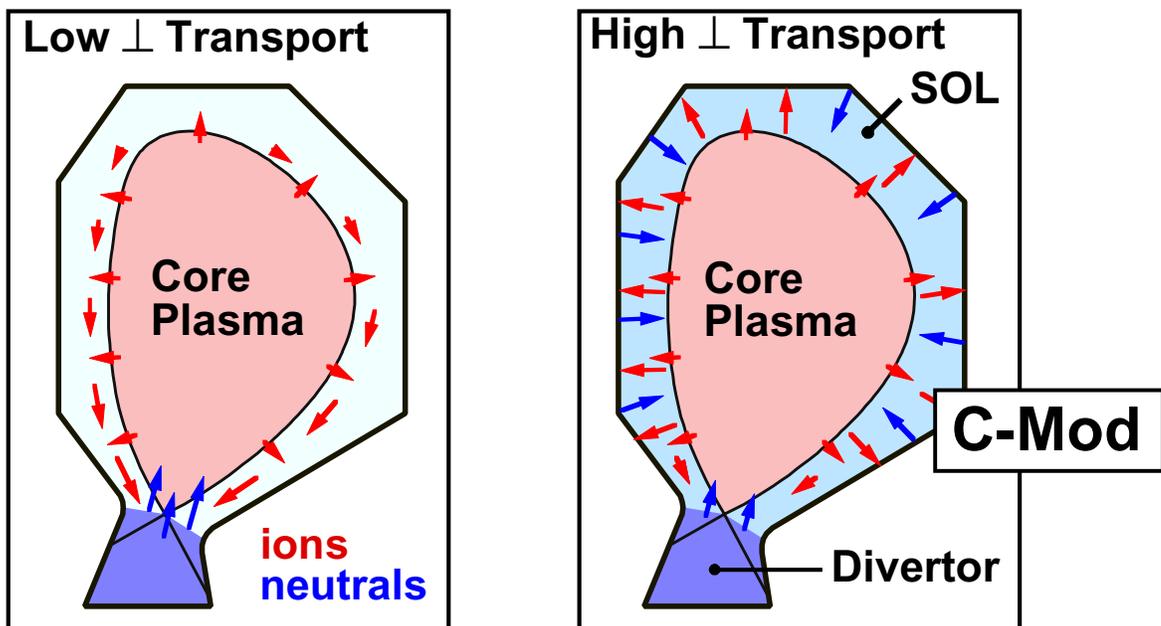
**B. LaBombard, R.L. Boivin, B. Carreras,
M. Greenwald, J. Hughes, B. Lipschultz,
D. Mossessian, C.S. Pitcher, J.L. Terry,
S.J. Zweben, Alcator C-Mod Group**

**Presented at the
International Atomic Energy Agency
Technical Committee Meeting
on DIVERTOR CONCEPTS
September 11 — 14, 2001, Aix-en-Provence, France**

Cross-field transport in SOL is a key physics issue...

- impacts tokamak operation and divertor design

- Γ_{\perp} determines level of plasma/wall interaction in main-chamber
 - neutral pressures (\Leftrightarrow confinement)
 - wall impurity sources
 - \Rightarrow impacts divertor design



Q: What will be operating regime in a reactor?

Cross-field transport in SOL is a key physics issue...

- must be understood for predictive modelling

- Heat convection across separatrix and SOL increases with plasma density
=> may precipitate divertor detachment!
- Particle transport is characterized by 'bursty', large-transport 'events'
=> diffusive transport paradigm in SOL/divertor simulations is inadequate!

- can directly impact core plasma

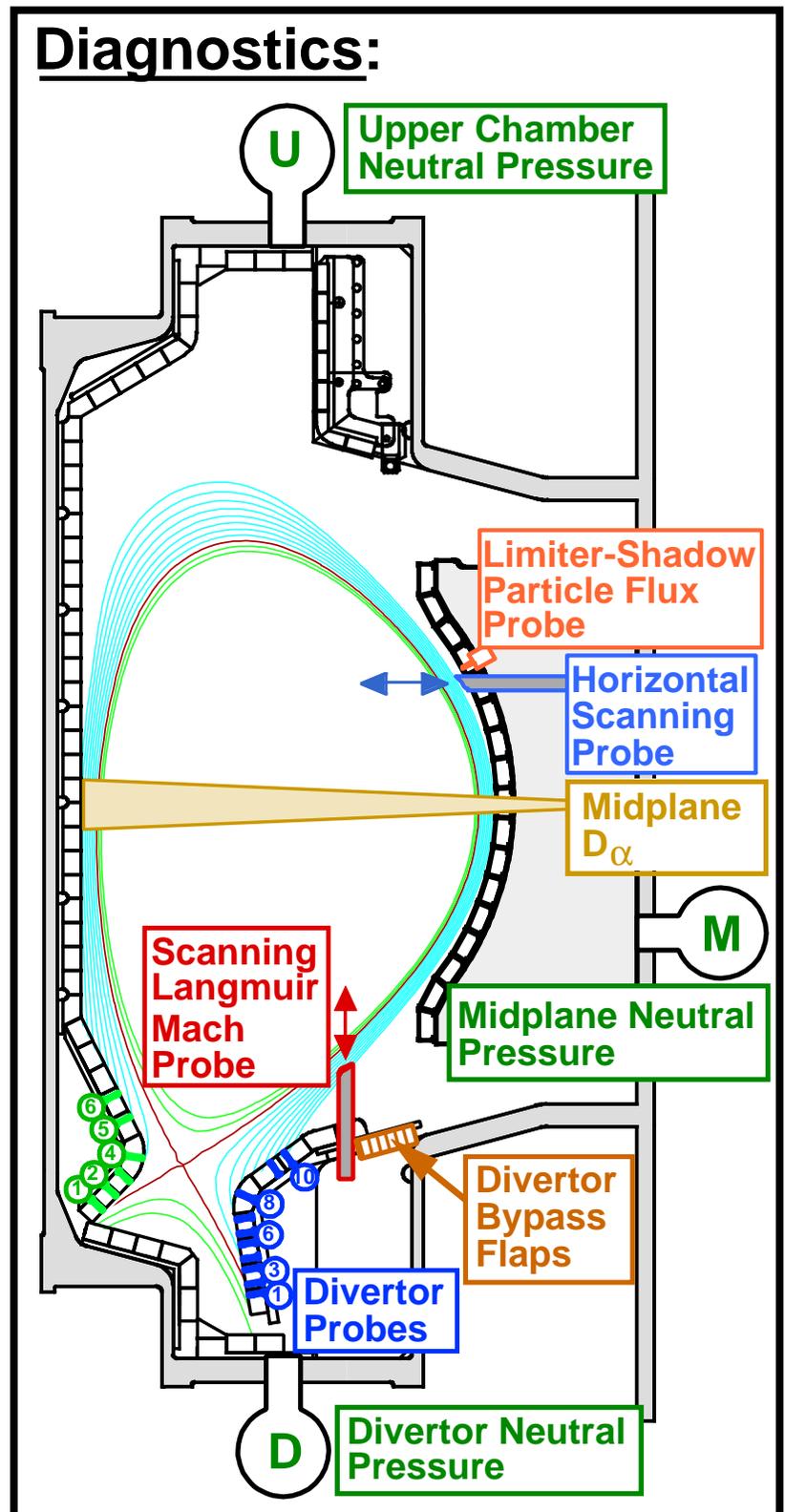
- Near the density limit, cross-field convection dominates heat fluxes through entire SOL
=> empirical scaling of tokamak density limit may be set by the physics of SOL transport!

=> Need to develop scaleable empirical and physics-based understandings of underlying transport physics

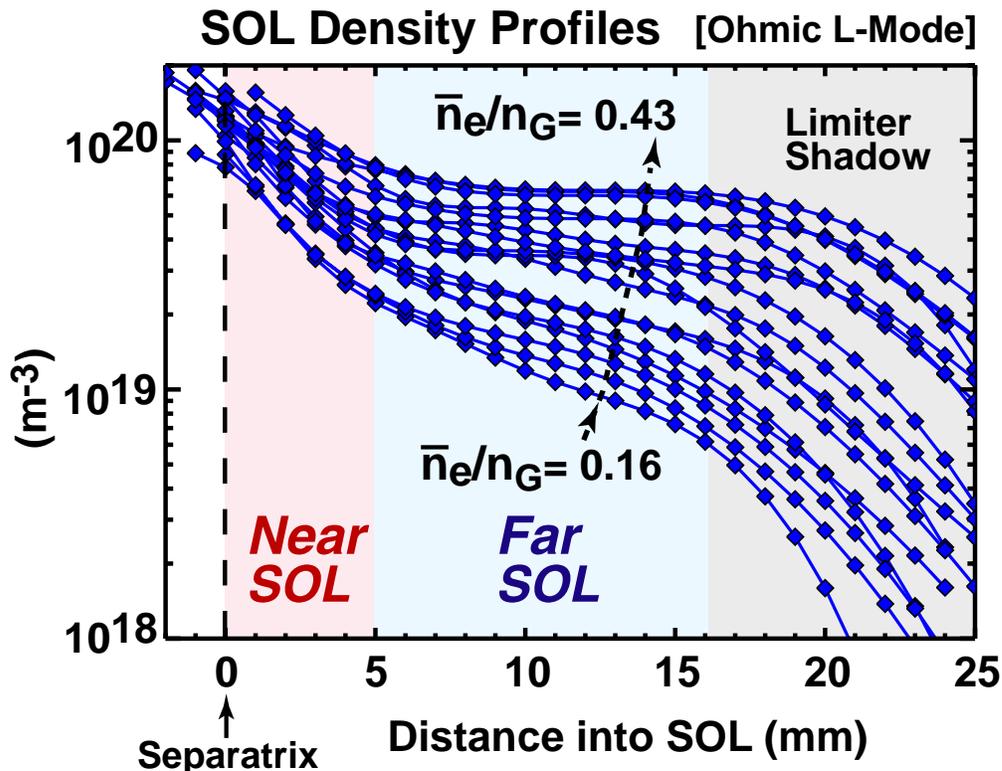
None exist at present!

Outline of Talk

- Main-Chamber Particle Balance
- Critical Cross-Field Flux for Particle Control
- Effective Cross-Field Particle Diffusivities (D_{eff}) & Scalings
- Cross-Field Heat Convection
- Character of SOL Fluctuations
- SOL Transport Physics and the Discharge Density Limit



Scrape-off Layer Density Profiles Exhibit a "Two-Exponential" Decay[†]



Near SOL: steep decay, $\lambda_n \sim 2$ to 8 mm

Far SOL: shallow decay, $\lambda_n \sim 8$ to > 100 mm

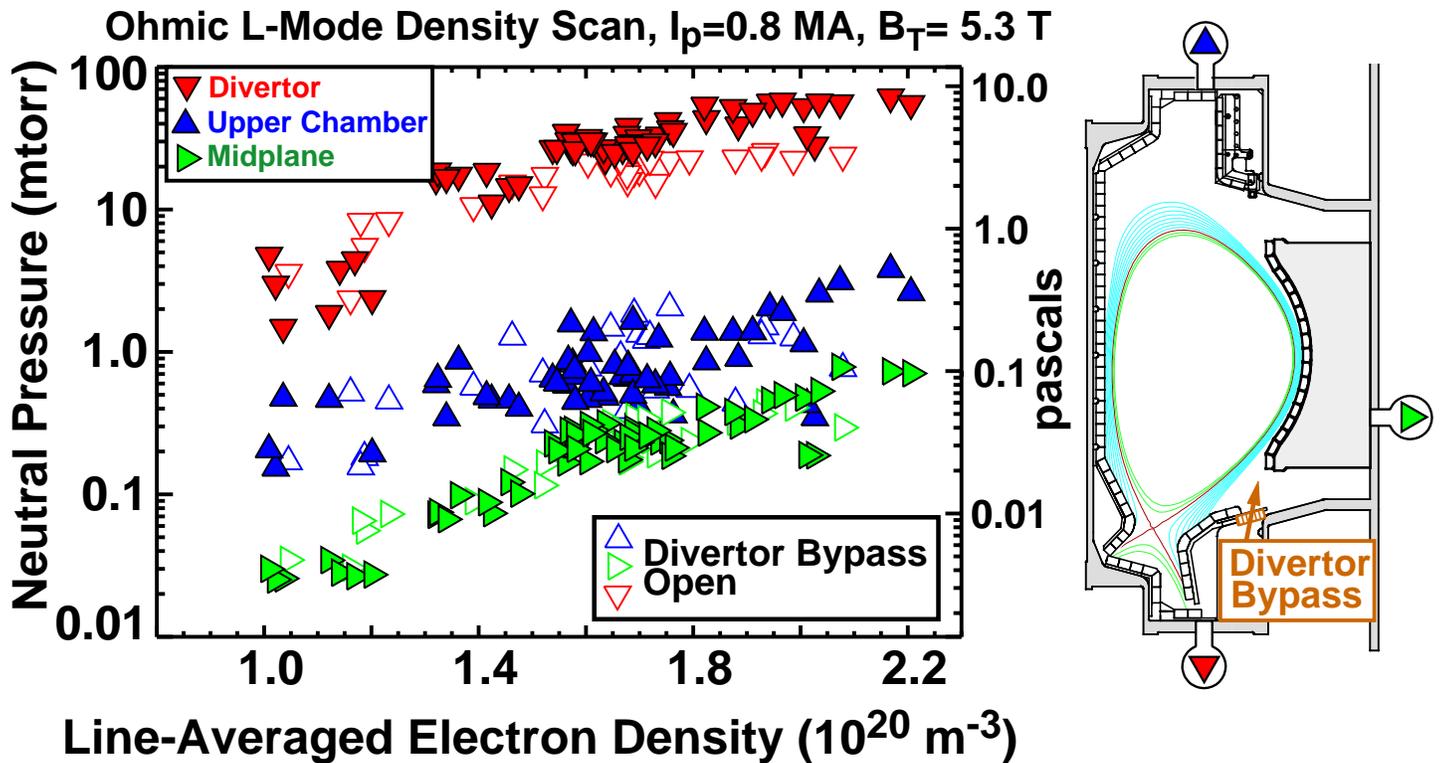
- At low \bar{n}_e , density at limiter edge is less than $\sim 1/10$ of separatrix density
- Density at limiter edge increases sharply with increasing \bar{n}_e

=> Always some level of main-chamber (limiter) recycling

Note: Similar *Far SOL* profiles are seen in H-mode discharges with the same midplane neutral pressure

[†] 'shoulders' on SOL profiles are prevalent in the literature: ASDEX, ASDEX-U, JT-60U, TEXT-U, ...

High Neutral Pressures Surround Core Plasma, Independent of Divertor Bypass



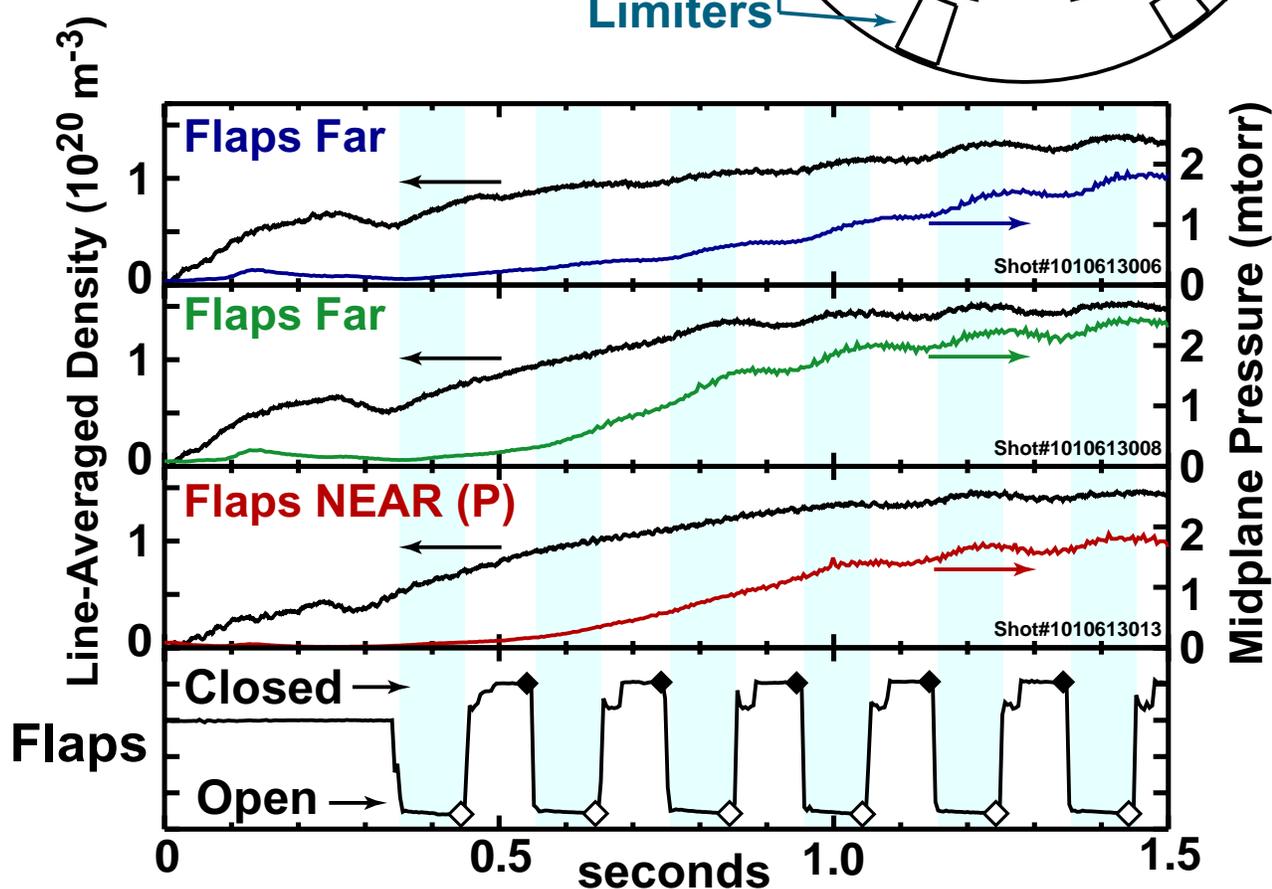
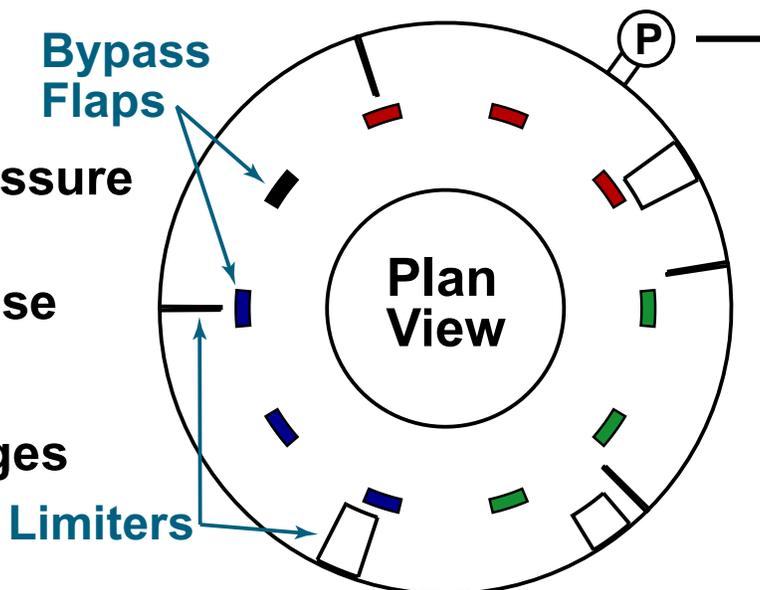
- **Divertor Pressure** drops by factor of 2 when bypass is opened at high density
 - Yet, robust relationship between **Midplane** pressure and line-averaged density is seen[†]
 - Neutral pressures in **Upper Chamber** can be higher than at **Midplane**!
- => implies large neutral fluxes (Γ_w) attack core plasma directly from main-chamber "wall" surfaces**

[†] as noted on ASDEX-U ...

Local Midplane Pressure is Insensitive to Neutral Leakage Through Local Divertor Bypass

Divertor Bypass Flap Experiment:

- Monitor midplane pressure at location (P)
- Dynamically open/close 3 flaps at different toroidal locations in 3 different discharges



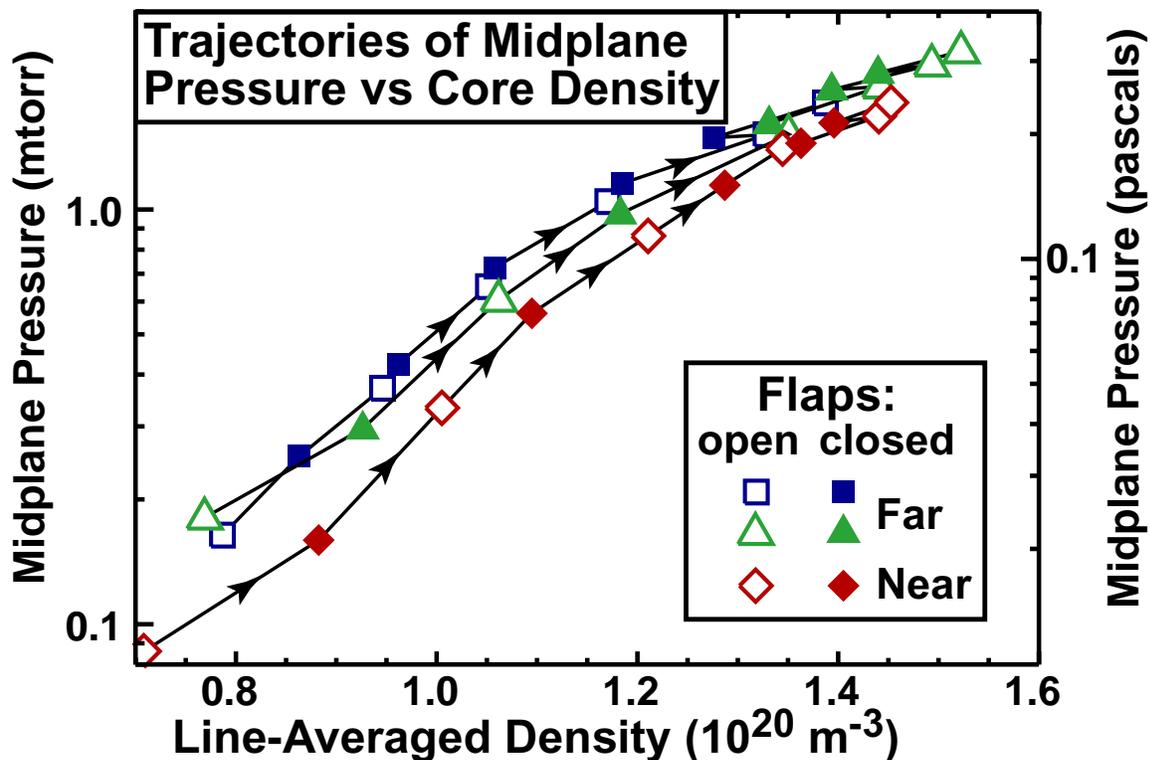
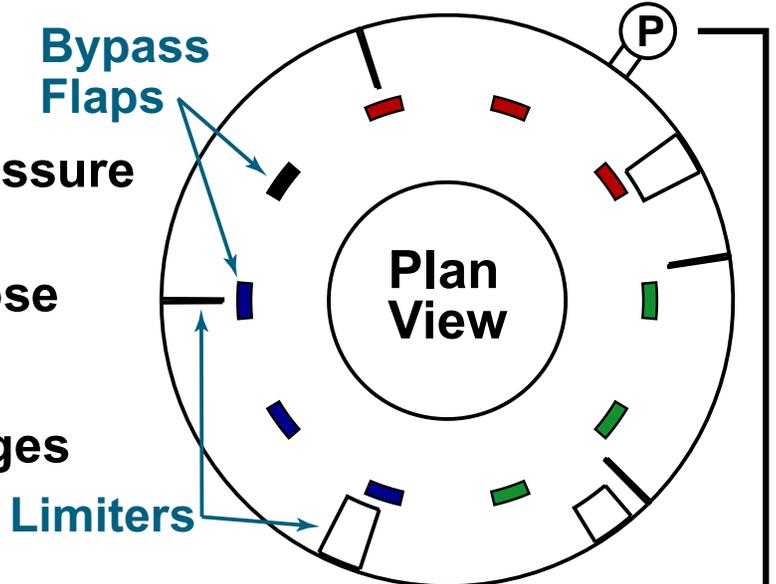
Local midplane pressure and line-averaged density perturbations are of similar magnitude, independent of toroidal location of active bypass flaps

=> Most of neutral leakage from divertor does not directly contribute to midplane pressure

Local Midplane Pressure Tracks Core Density, Independent of Local Divertor Bypass State

Divertor Bypass Flap Experiment:

- Monitor midplane pressure at location (P)
- Dynamically open/close 3 flaps at different toroidal locations in 3 different discharges

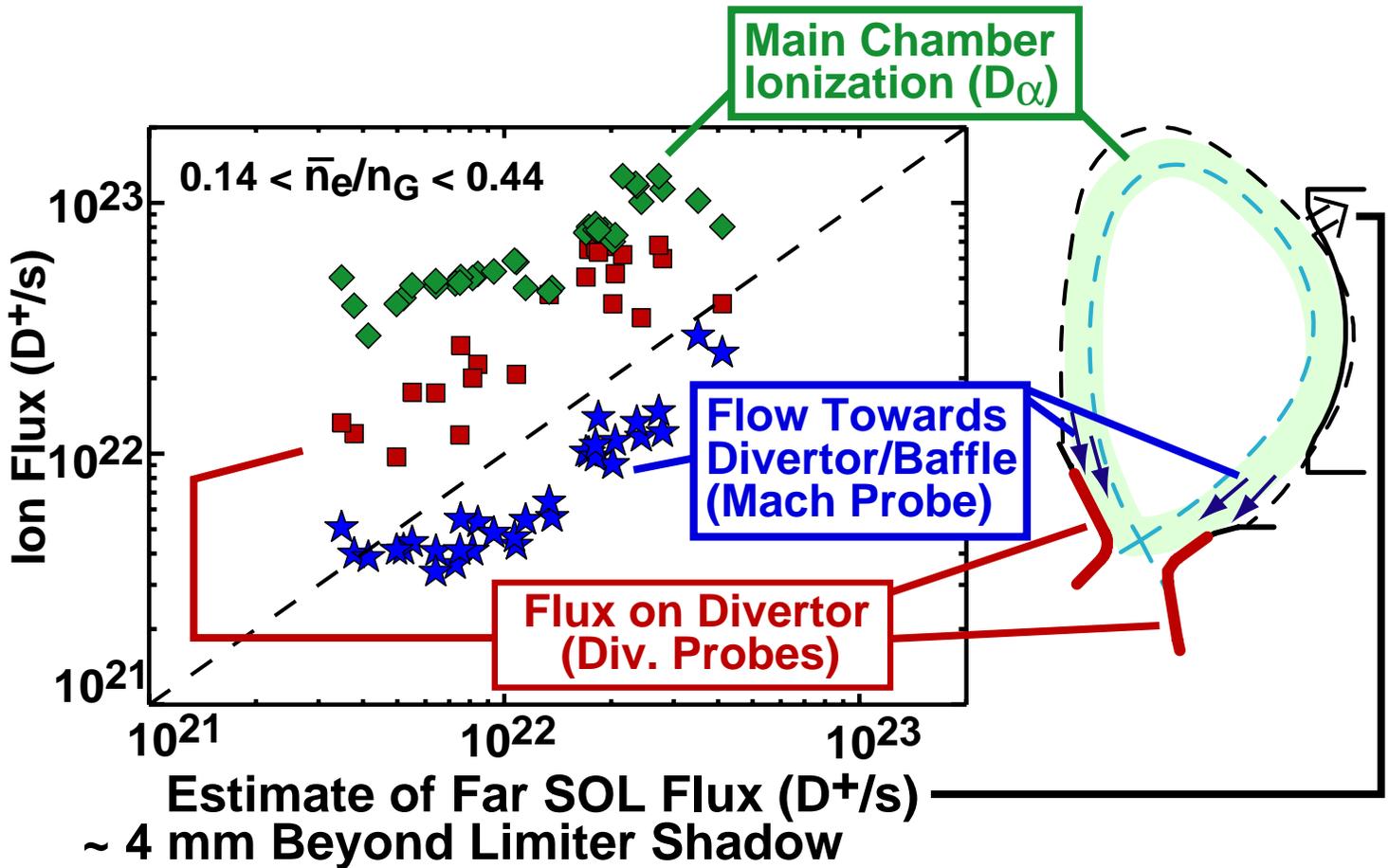


Local midplane pressure tracks line-averaged density, independent of toroidal location of active bypass flaps

=> Midplane neutral pressure is set primarily by recycling on main-chamber surfaces rather than divertor leakage[†]

[†]see Lipschultz et al., this conference

Main Chamber Ionization and Ion Fluxes to Main Chamber Limiters are Large Compared to Flows Towards Div./Baffle



- Recycling in **Main Chamber SOL** is primarily balance by fluxes onto main-chamber walls[†]
- **Poloidal flows** to divertor/baffle are weak

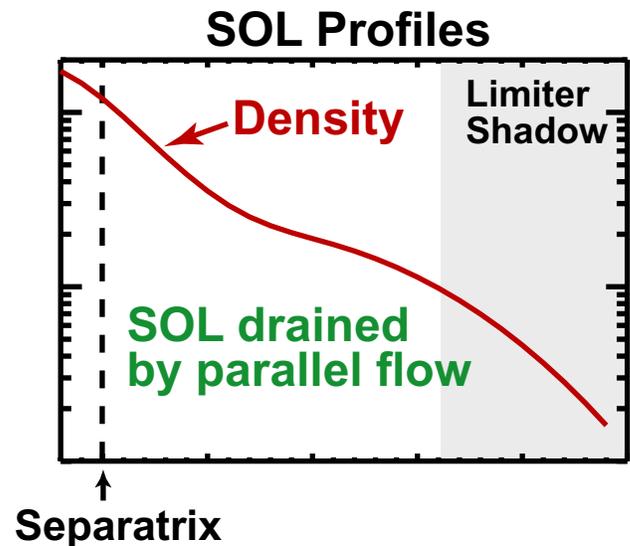
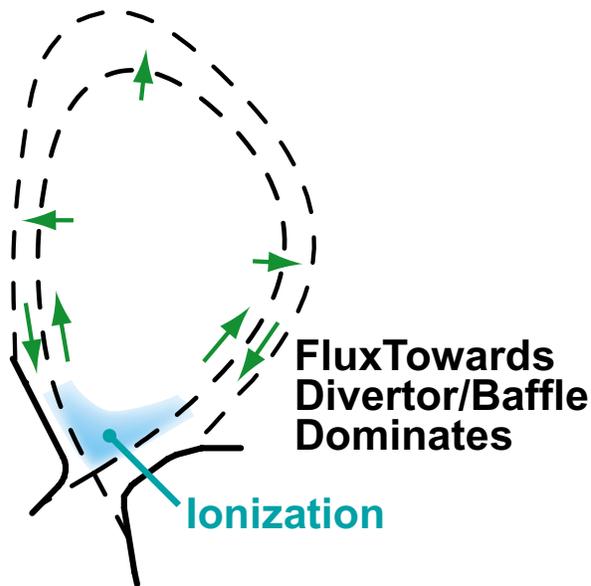
Main-Chamber Recycling Regime (MCR) persists over wide parameter range

[†]M.V. Umansky, et al. Phys. Plasmas 5, 3373 (1998).

A New View of Particle Transport Processes in SOL

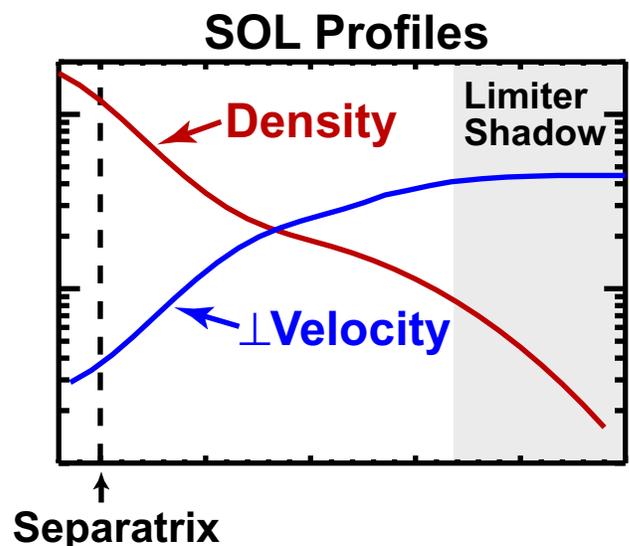
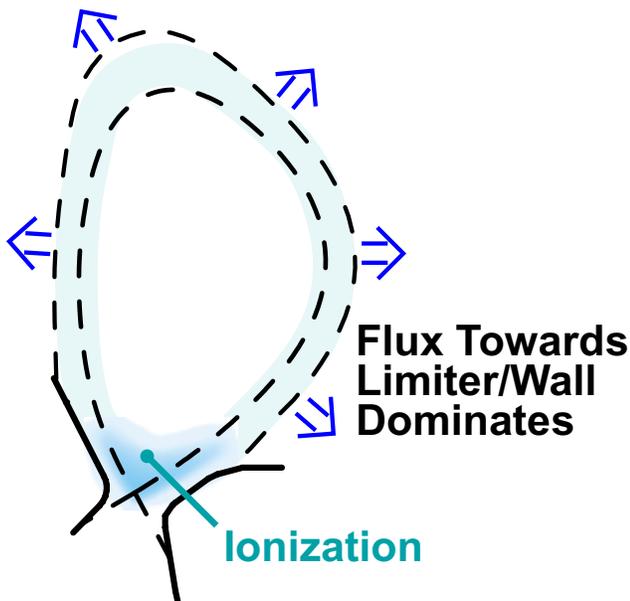
Old Paradigm:

- **SOL density decays "exponentially"** because...
plasma drains along field lines towards divertor/baffle



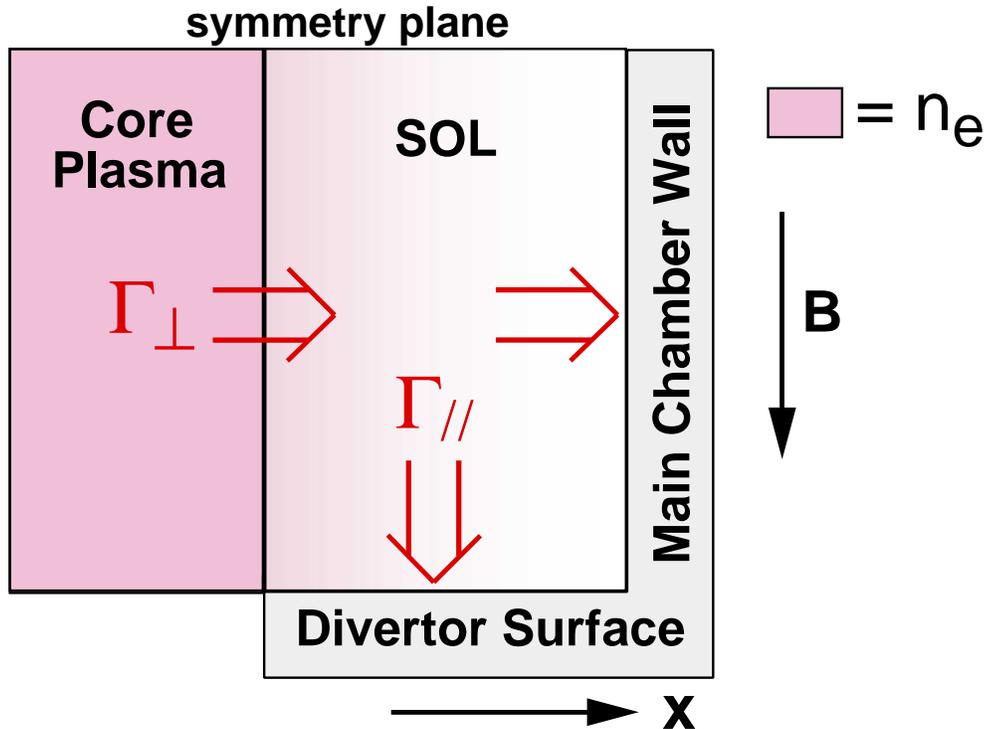
New Paradigm:

- **SOL density decays "exponentially"** because...
cross-field transport velocity increases across SOL,
maintaining cross-field flux towards wall



Q: What level of cross-field transport leads to loss of particle control in the main chamber?

Consider simplified SOL:



- Plasma Continuity

$$\nabla \cdot \underline{\Gamma} = n_0 n_e k_{\text{ion}}$$

- Flux Surface Average $\langle \dots \rangle$

$$\frac{\partial}{\partial x} \langle \Gamma_{\perp} \rangle \approx n_e \left[\langle n_0 \rangle k_{\text{ion}} - \frac{C_s}{2L} \right]$$

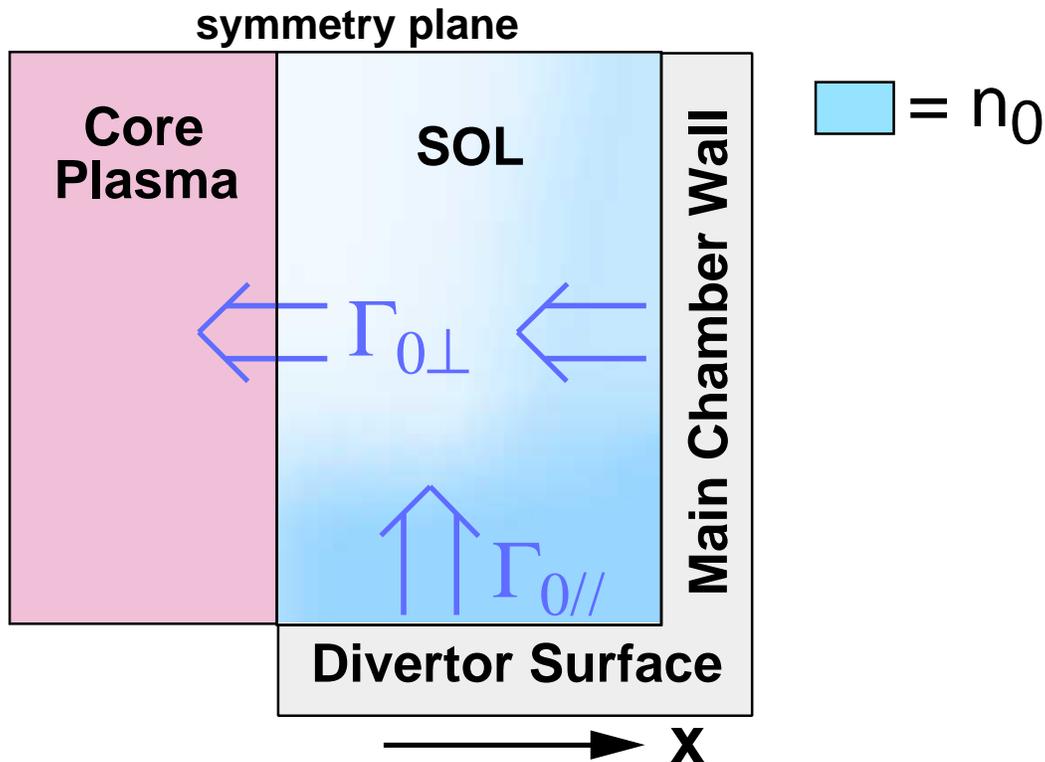
Note #1: If a critical flux-surface averaged neutral density is exceeded, $\langle \Gamma_{\perp} \rangle$ must increase across SOL[†]

$$\langle n_0 \rangle_{\text{crit}} \sim \frac{C_s}{2k_{\text{ion}} L}$$

[†]a well known over-ionization condition (e.g., Stangeby et al.)

Q: What level of cross-field transport leads to loss of particle control in the main chamber?

Note #2: Mass balance requires: $\langle \Gamma_{0\perp} \rangle = - \langle \Gamma_{\perp} \rangle$



Note #3: Neutral Flux requires minimum Neutral Density:

$$\langle n_0 \rangle \geq \sim - \langle \Gamma_{0\perp} \rangle / v_{th0\perp} \quad (v_{th0\perp} \sim \text{set by CX})$$

A: If $\langle \Gamma_{\perp} \rangle$ exceeds $\sim \frac{C_s v_{th0\perp}}{2k_{ion} L}$ at $x=x_0$ then $\langle \Gamma_{\perp} \rangle$ increases with x for $x > x_0$ †

=> Result is insensitive to separatrix-wall distance and divertor geometry!

=> Main-chamber particle control depends critically on the level of \perp particle transport

† 'critical flux' is comparable to level of $\langle \Gamma_{\perp} \rangle$ observed in C-Mod

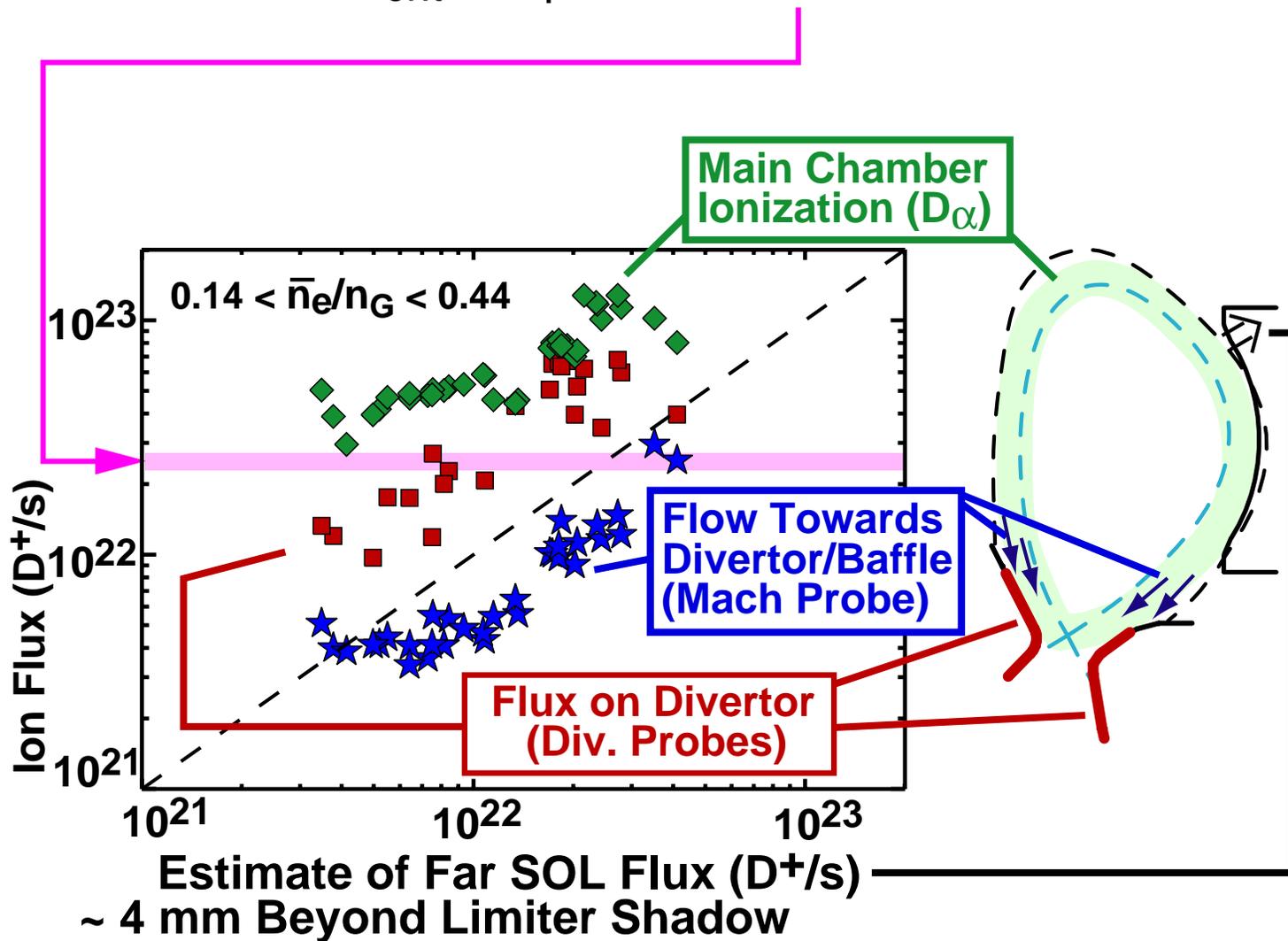
Main Chamber Ionization (D_α) in C-Mod Exceeds Upper Bound Estimate of $\langle \Gamma_\perp \rangle_{\text{crit}}$

Estimate with $T_0 \approx T_i \approx T_e$, free-streaming neutrals

$$\langle \Gamma_\perp \rangle_{\text{crit}} \approx 2 \times 10^{20} \frac{T \text{ (eV)}}{q R \text{ (m)}} \text{ (m}^{-2} \text{ s}^{-1}\text{)}$$

Typical Number (with $T \sim 50$ eV):

$$\langle \Gamma_\perp \rangle_{\text{crit}} A_{\text{sep}} \sim 2.6 \times 10^{22} \text{ (s}^{-1}\text{)}$$



Outline of Talk

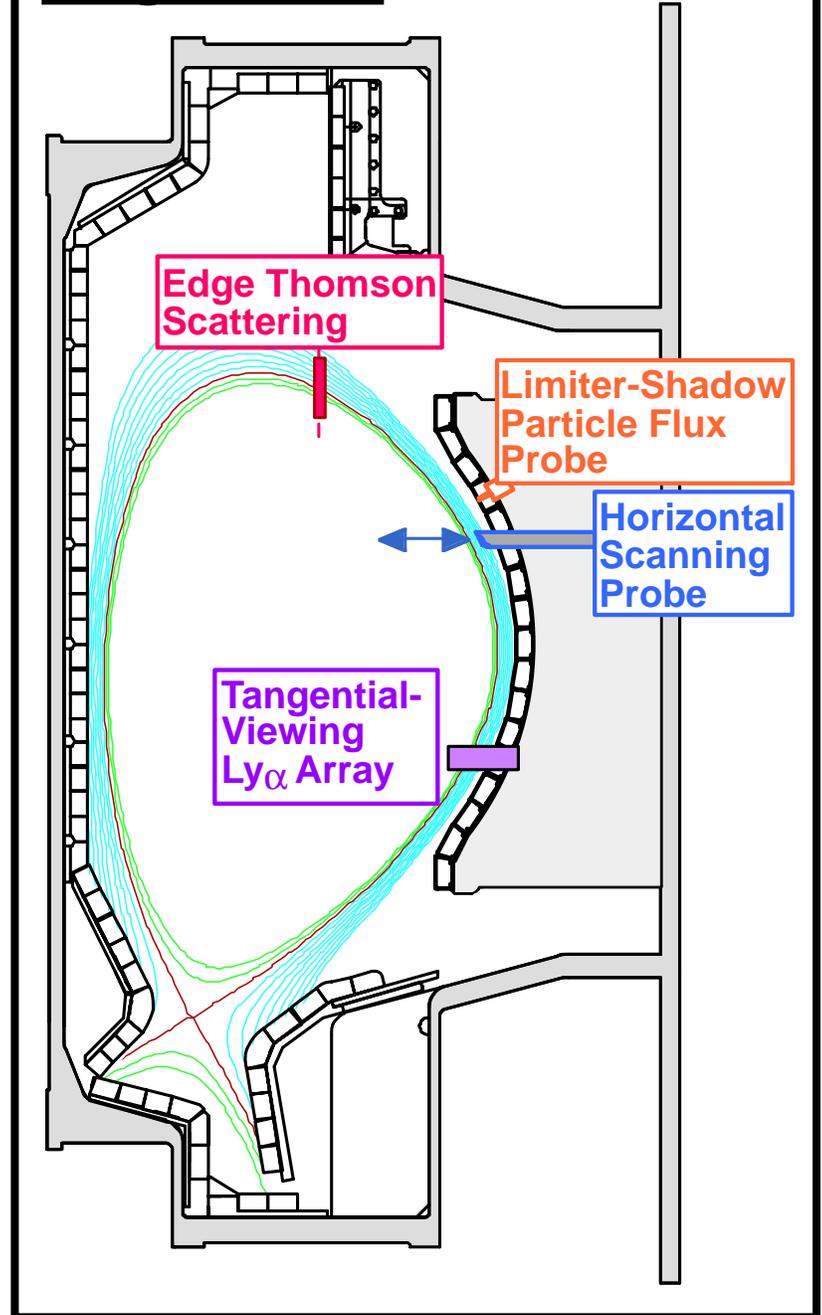
- Main-Chamber Particle Balance
- Critical Cross-Field Flux for Particle Control

• **Effective Cross-Field Particle Diffusivities (D_{eff}) & Scalings**

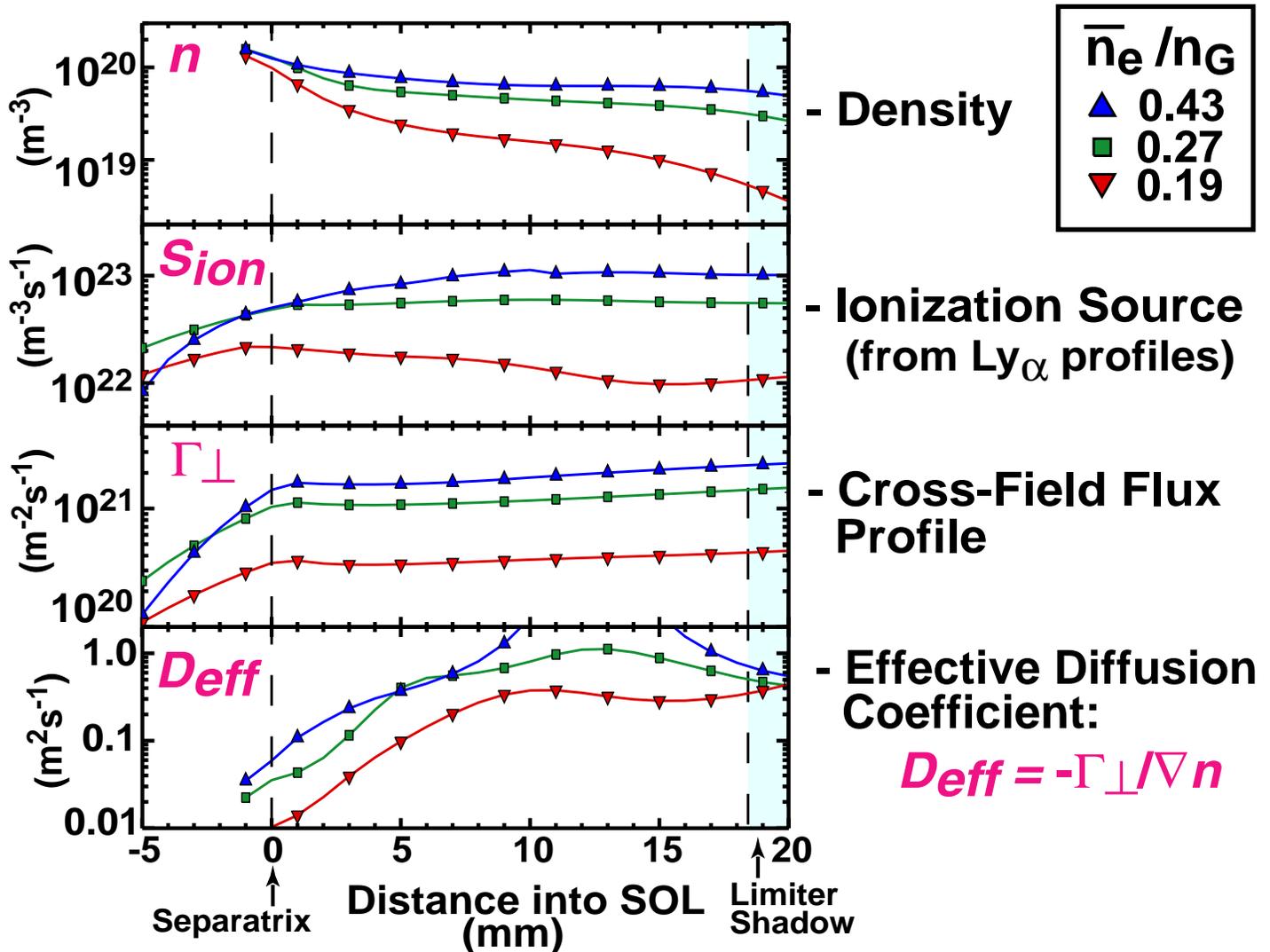
• **Cross-Field Heat Convection**

- Character of SOL Fluctuations
- SOL Transport Physics and the Discharge Density Limit

Diagnostics:



In MCR Regime, Cross-Field Diffusion Coefficient Profiles (D_{eff}) can be Inferred Directly from Profile Measurements[†]



- Persistent trend of D_{eff} increasing by ~ 10 or more with distance from separatrix^{††}
 \Rightarrow variation in D_{eff} reflects variation in ∇n
- D_{eff} increases with discharge density
 $\Rightarrow \Gamma_\perp$ gets larger, ∇n gets smaller

[†]Method benchmarked against UEDGE modeling

^{††} $D_{eff}(\chi_{eff})$ increasing seen before: ASDEX, JT-60, JET, ...

Magnitude of D_{eff} in Near SOL is Correlated with Collisionality in Near SOL

64 Ohmic L-Mode Datapoints:

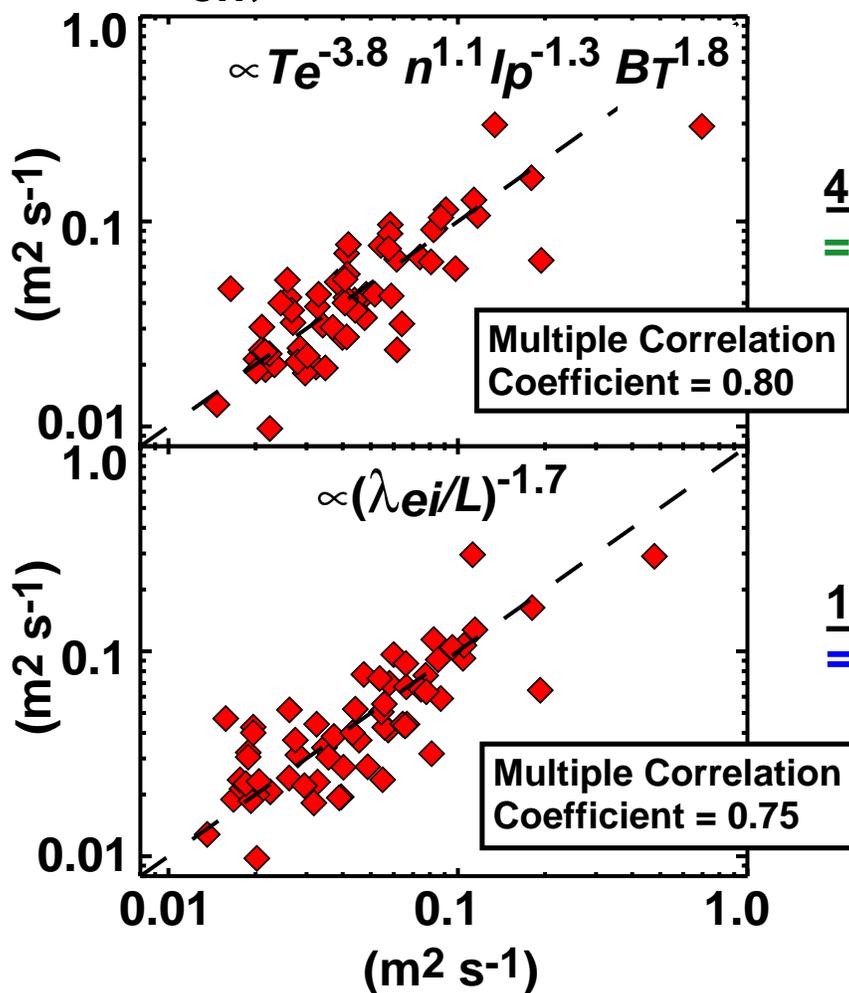
$$0.8 < \bar{n}_e < 2.5 \times 10^{20} \text{ m}^{-3}$$

$$0.6 < I_p < 1.0 \text{ MA}$$

$$4 < B_T < 6 \text{ tesla}$$

$$0.14 < \bar{n}_e/n_G < 0.47$$

Regression Analysis of D_{eff} , 2 mm into SOL



4 Parameter Regression:

=> Suggests (B_T/I_p) , q , or L dependence

1 Parameter Regression:

=> Statistics point to (λ_{ei}/L) as most relevant parameter

=> D_{eff} correlates with local collisionality:

$$D_{eff} \sim (\lambda_{ei}/L)^{-1.7}$$

Trend: $n_e/n_G \uparrow \Rightarrow \lambda_{ei}/L \downarrow \Rightarrow D_{eff} \uparrow$ near sep.

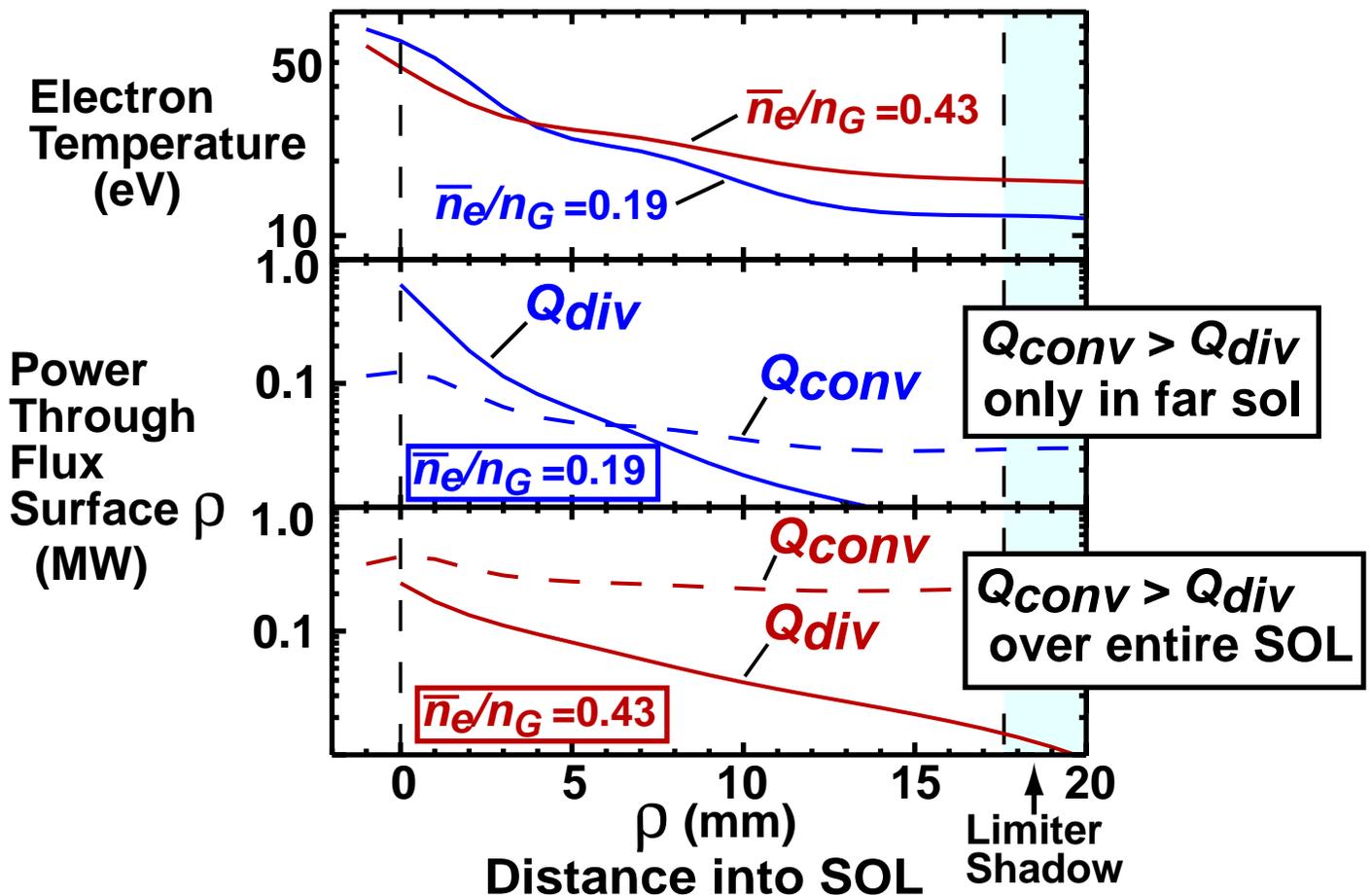
Cross-Field Heat Convection to Limiter/Wall Competes with Parallel Conduction Losses to Divertor at Moderate \bar{n}_e/n_G

Finite T_e on open field lines

=> power conducted to divertor: $Q_{div} \propto \int_{\rho}^{\infty} T^{7/2}/L \delta\rho'$

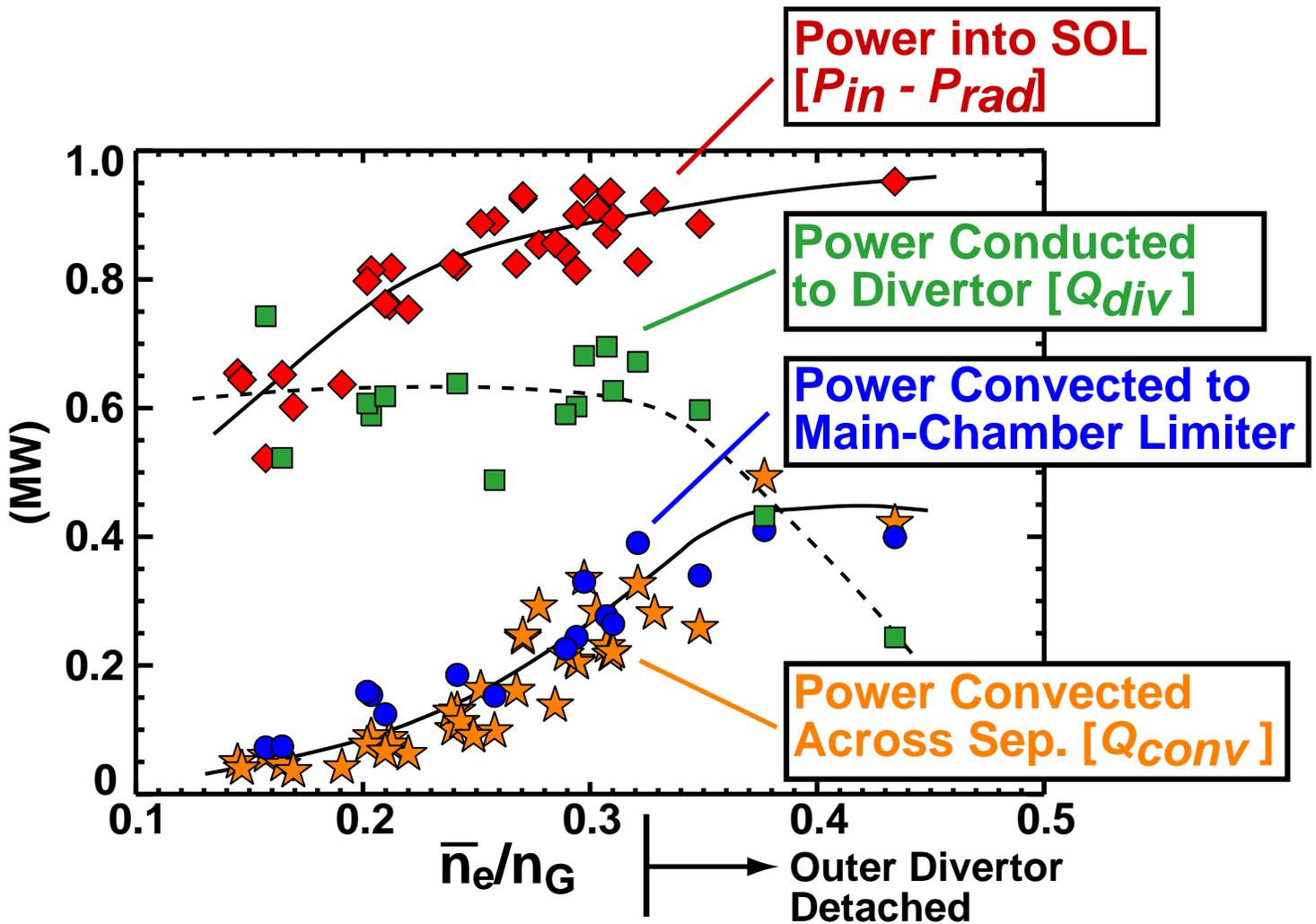
Cross-field particle fluxes (Γ_{\perp})

=> power convected: $Q_{conv} \sim 5 T_e \Gamma_{\perp} A_{sep}$



- At low density, heat losses in *Near SOL* are dominated by parallel conduction to Divertor
- At moderate density, cross-field heat convection to Limiter/Wall exceeds conduction losses to Divertor/Baffle over *entire SOL*

Cross-Field Convection Increases with \bar{n}_e/n_G , Affecting SOL Power Balance



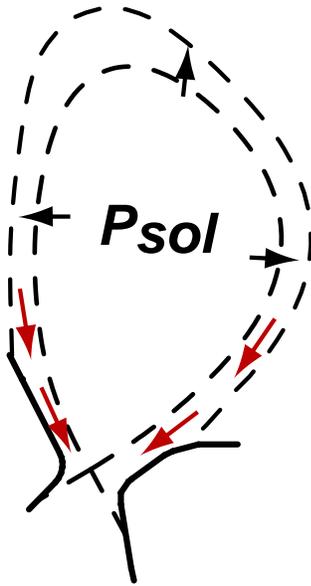
- At low density, parallel conduction to Divertor dominates SOL power balance
 - At moderate density, cross-field heat convection to Limiter/Wall becomes important
- ⇒ Cross-field convection losses to main-chamber wall may precipitate divertor detachment

A New View of Heat Transport Processes in SOL

(in absence of a "radiating mantle")

Old Paradigm:

- **Parallel e⁻ conduction** to divertor dominates heat losses in Near SOL region



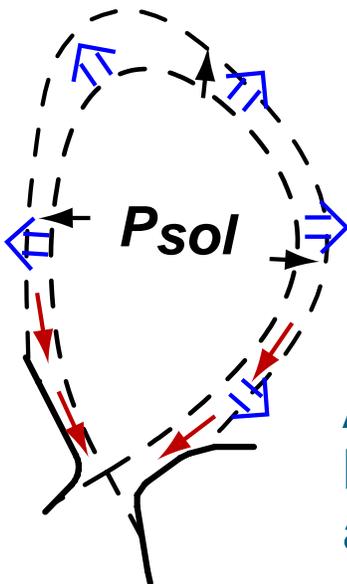
T_e at separatrix (T_{sep}) is a weak function of λ_{Te} and SOL power (P_{sol}):

$$T_{sep} \propto (P_{sol} / \lambda_{Te})^{2/7}$$

Modified Paradigm:

- **Parallel e⁻ conduction** and **cross-field heat convection** contribute to heat losses in Near SOL region

(CX is typically a minor player)



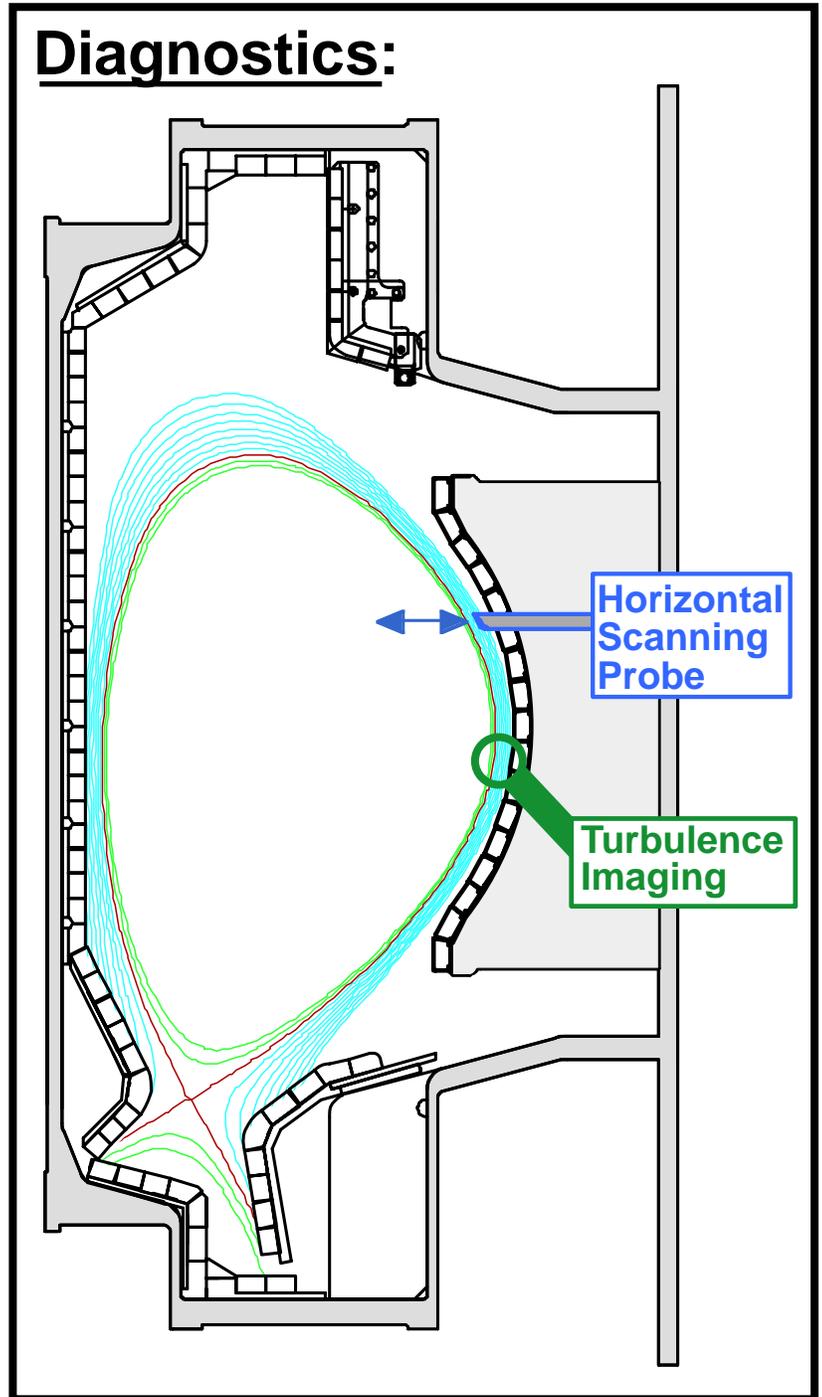
At low collisionality, parallel conduction regulates T_{sep} :

$$T_{sep} \propto (P_{sol} / \lambda_{Te})^{2/7}$$

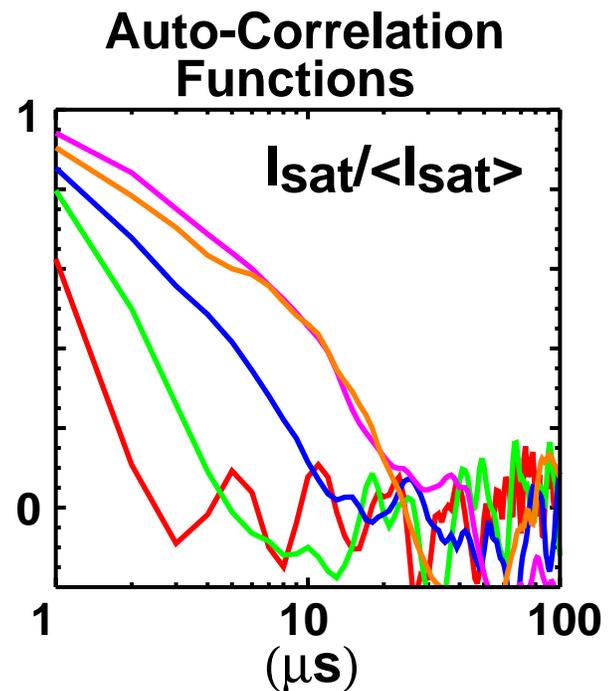
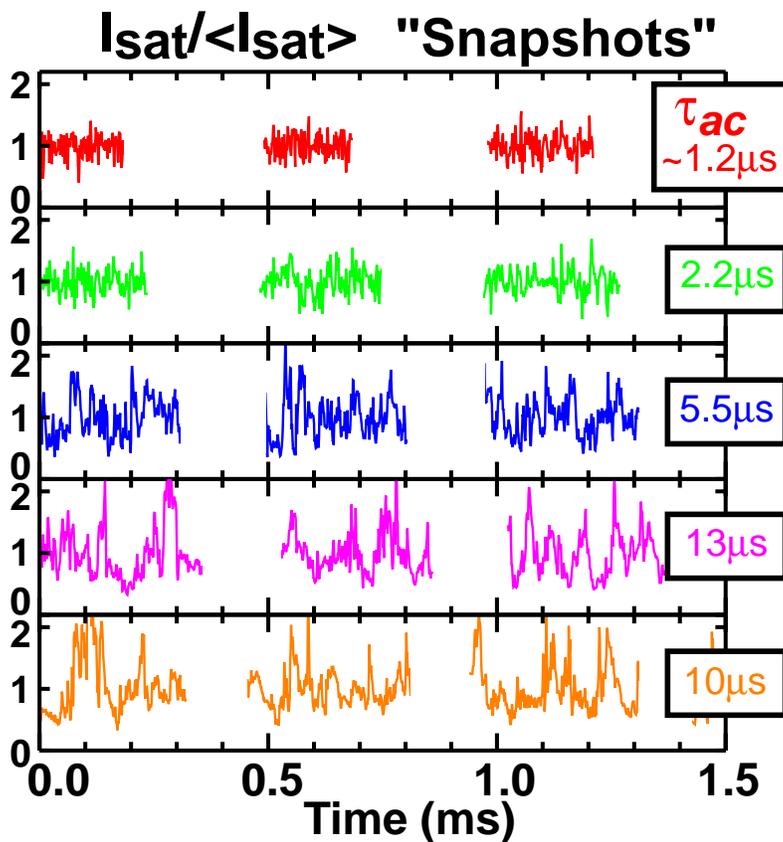
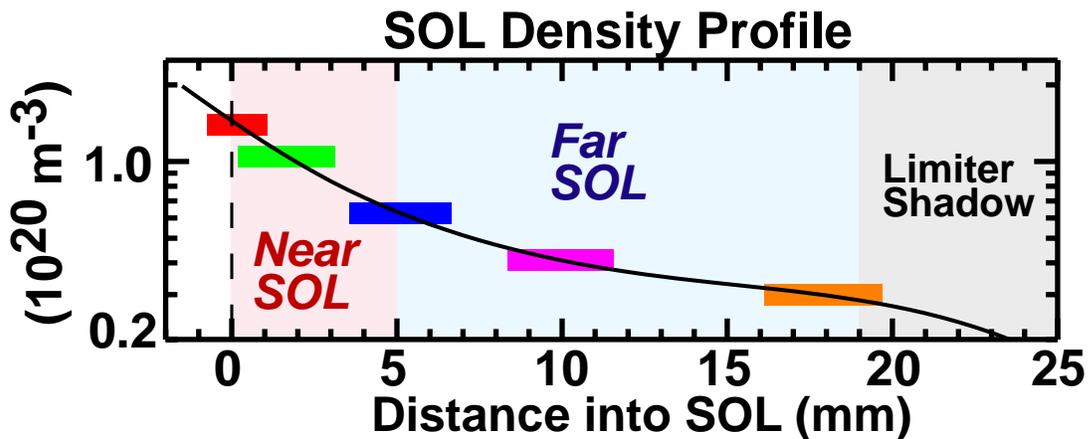
At high collisionality, heat convection becomes large, T_{sep} is reduced and is no longer "regulated" by this law!

Outline of Talk

- Main-Chamber Particle Balance
- Critical Cross-Field Flux for Particle Control
- Effective Cross-Field Particle Diffusivities (D_{eff}) & Scalings
- Cross-Field Heat Convection
- **Character of SOL Fluctuations**
- SOL Transport Physics and the Discharge Density Limit



Fluctuations Exhibit Different Character in *Near* and *Far* SOL Regions



Near SOL (steep n profile):

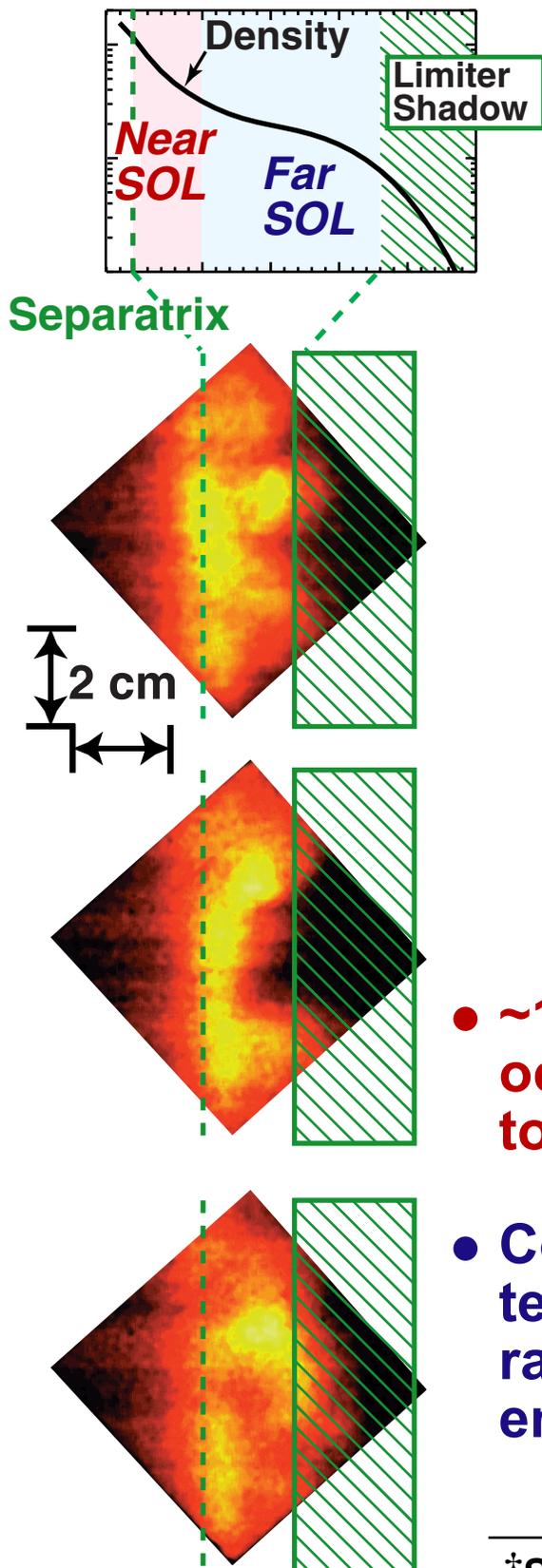
-> moderate amplitude, "random" fluctuations

Far SOL (flatter n profile):

-> large amplitude, intermittent I_{sat} "bursts"

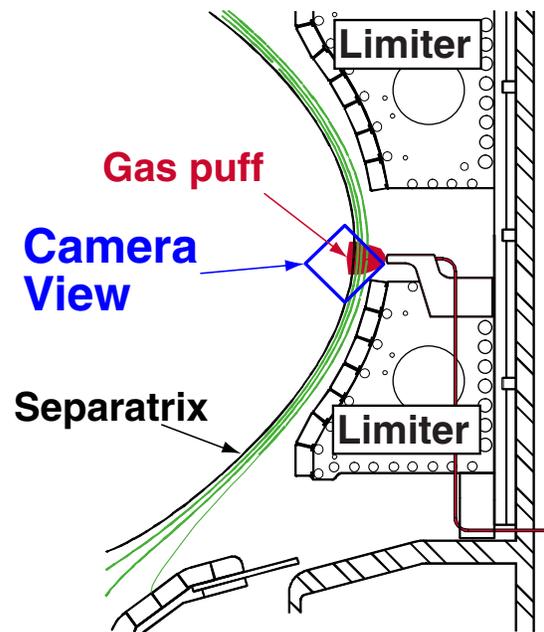
=> Consistent with $D_{\text{eff}} \uparrow$ with distance into SOL

2-D Turbulence Imaging: Intermittent, ~1 cm Scale "Blobs" of Emission Extend into Far SOL



Turbulence Imaging: †

Camera looks along field lines at a D_2 gas puff



2 s exposure times
17 ms between exposures

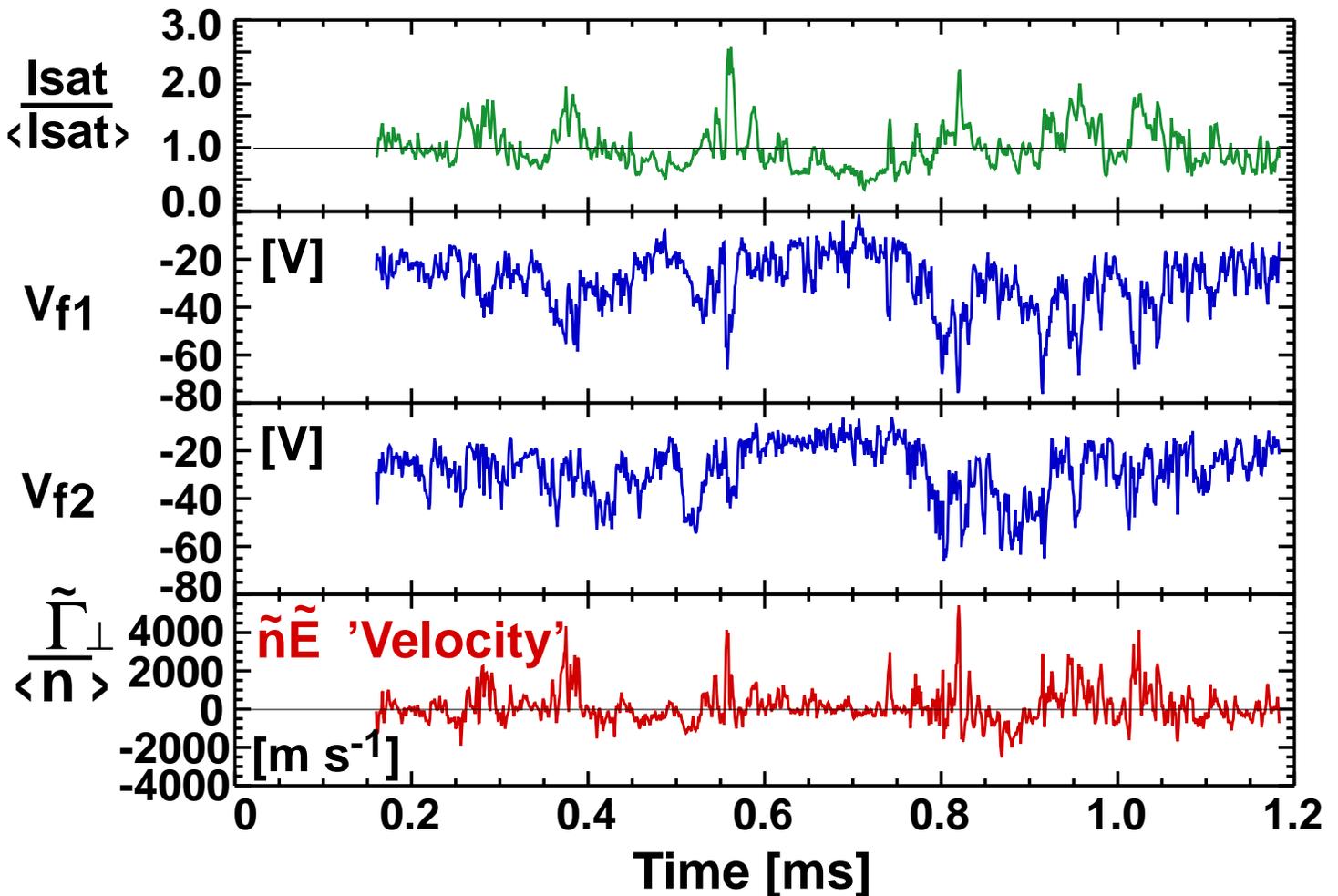
- ~1 cm scale blobs intermittently occupy Far SOL zone, and extend to Limiter Shadow
- Consistent with large density and temperature (?) perturbations rapidly transporting particles and energy to Limiter/Walls

†S. Zweben, J.L. Terry, R. Maqueda

Time-History of Fluctuation-Driven Flux Shows Large-Amplitude, Bursty Behavior

Data from Midplane Probe,
7 mm outside LCFS

Particle flux estimate neglects \tilde{T}_e



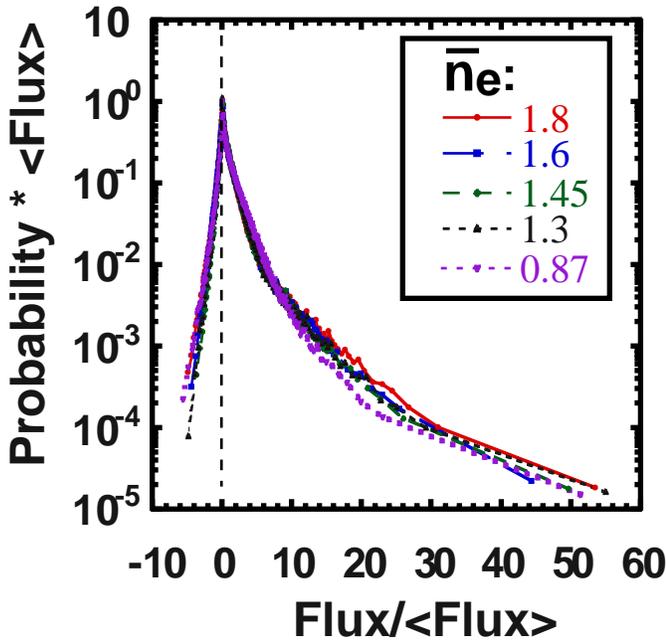
Time-averaged transport velocity: $\frac{\langle \tilde{\Gamma}_{\perp} \rangle}{\langle n \rangle} \sim 120 \text{ m s}^{-1}$

'Bursts' in transport velocity exceed 2000 m s^{-1}

What is influence of \tilde{T}_e on these estimates?

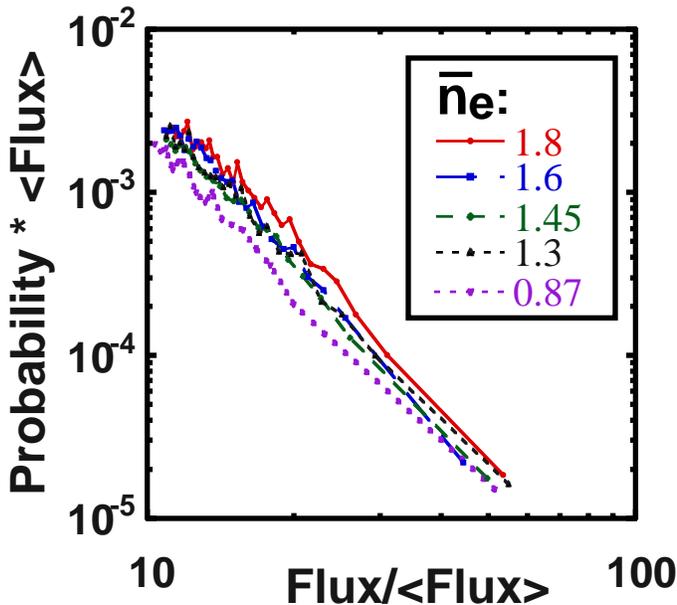
PDF of Fluctuation-Driven Particle Flux Exhibits Power-Law Tail, Independent of \bar{n}_e [†]

Γ_{\perp} Inferred from Midplane Probe, ~7 mm outside LCFS



Similar $P(\Gamma_{\perp})$ for all densities
=> change in $\langle\Gamma_{\perp}\rangle$ not associated with change in $P(\Gamma_{\perp})$

Always has SOC-Like behavior: ††



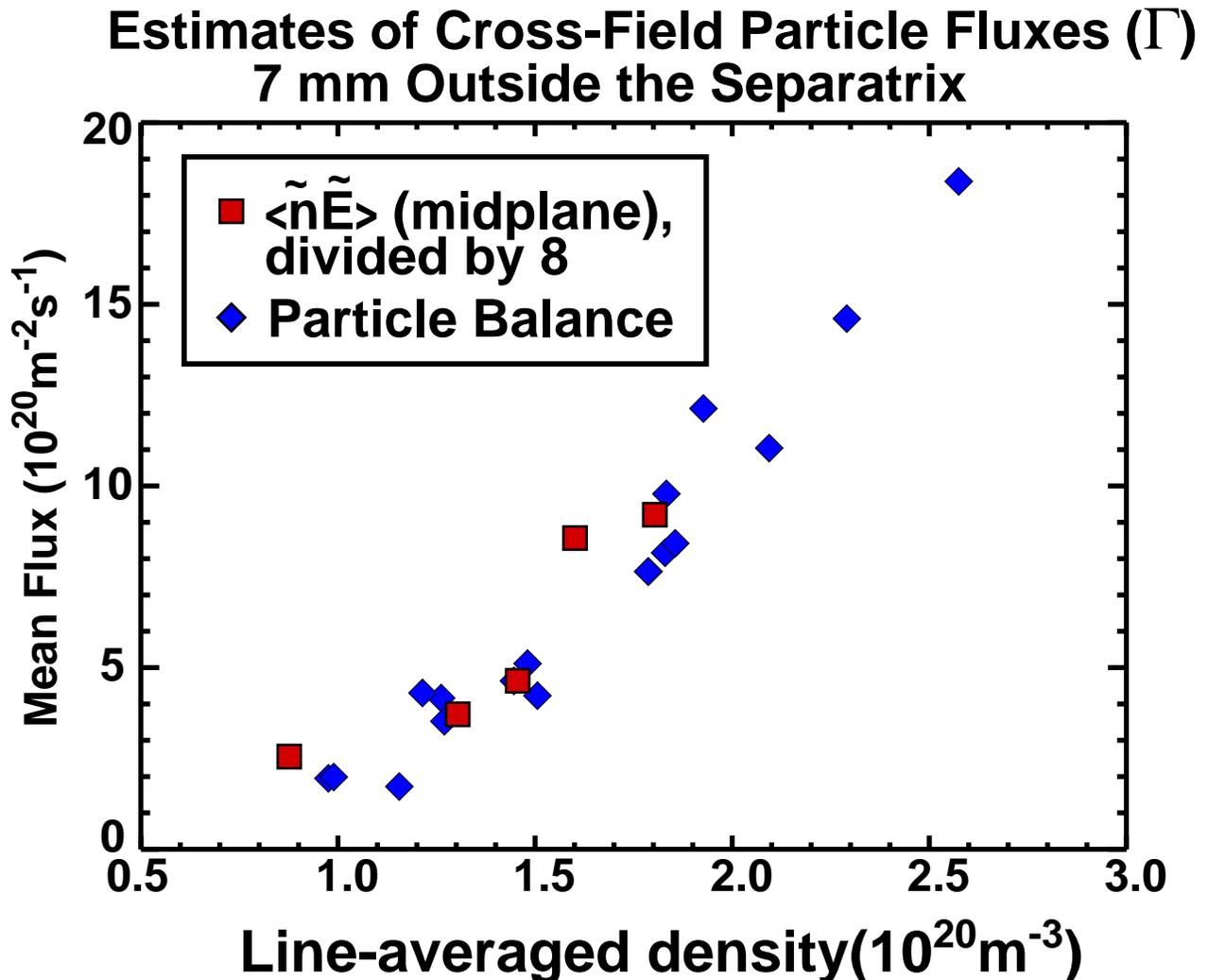
Positive Γ_{\perp} 'events' with Γ_{\perp} greater than $4.6 \langle\Gamma_{\perp}\rangle$ account for 50% of total particle transport.

These events happen ~5% of time.

[†]Analysis by B. A. Carreras, V. E. Lynch.

^{††}Bursty, SOC-like behavior of SOL plasma is universally seen in SOL plasmas including non-tokamak devices.

$\langle \tilde{n}\tilde{E} \rangle$ -Derived Cross-Field Flux Shows Similar Trend with \bar{n}_e as Particle Balance-Derived Flux



$\langle \tilde{n}\tilde{E} \rangle$ - inferred Γ does not account for possible \tilde{T}_e

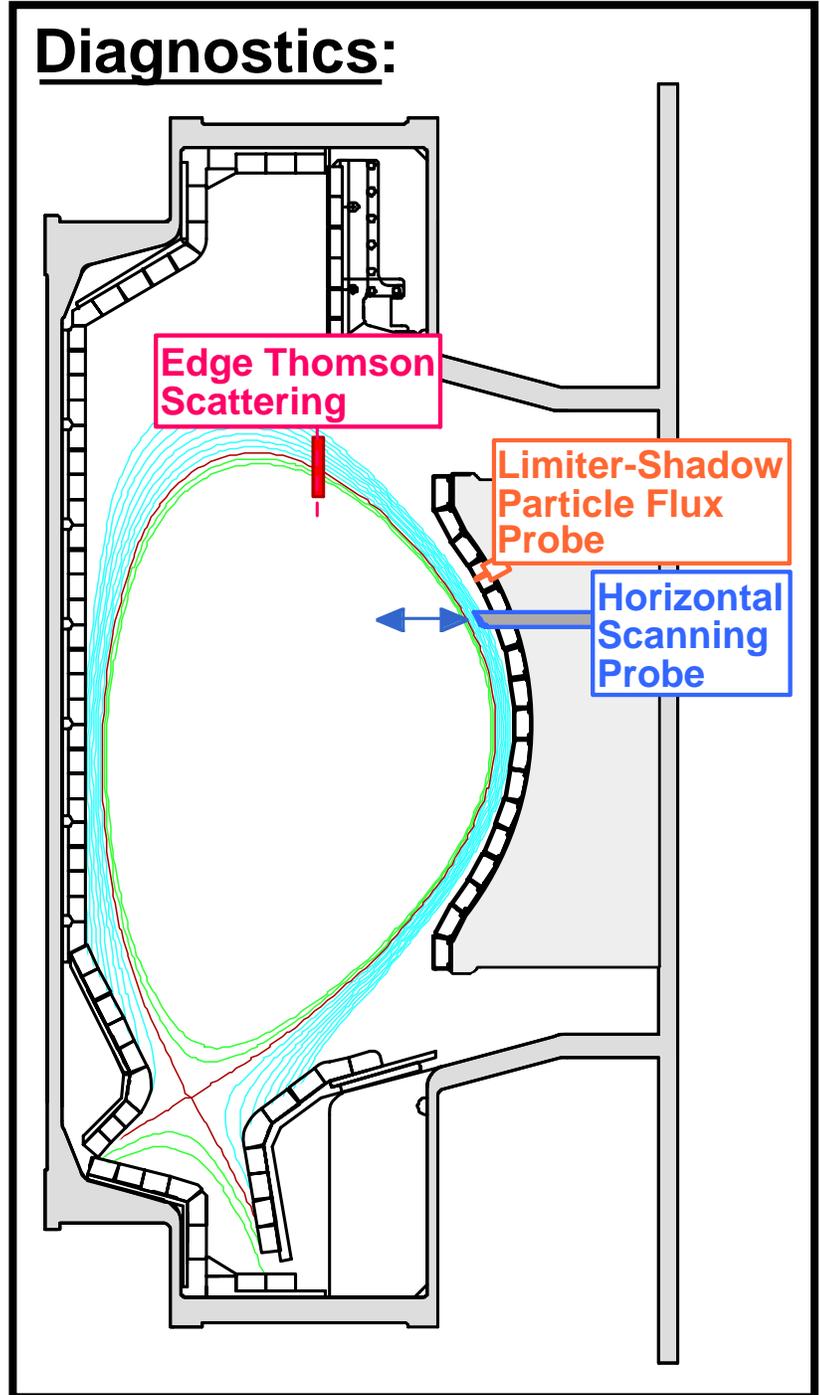
- Γ from both methods show similar trend, nonlinearly increasing with \bar{n}_e [†]
- Magnitude of Γ inferred from Midplane Probe $\langle \tilde{n}\tilde{E} \rangle$ is a factor of ~8 times larger than that derived from particle balance

[†] $\langle \tilde{n}\tilde{E} \rangle$ increasing with density seen before: ASDEX

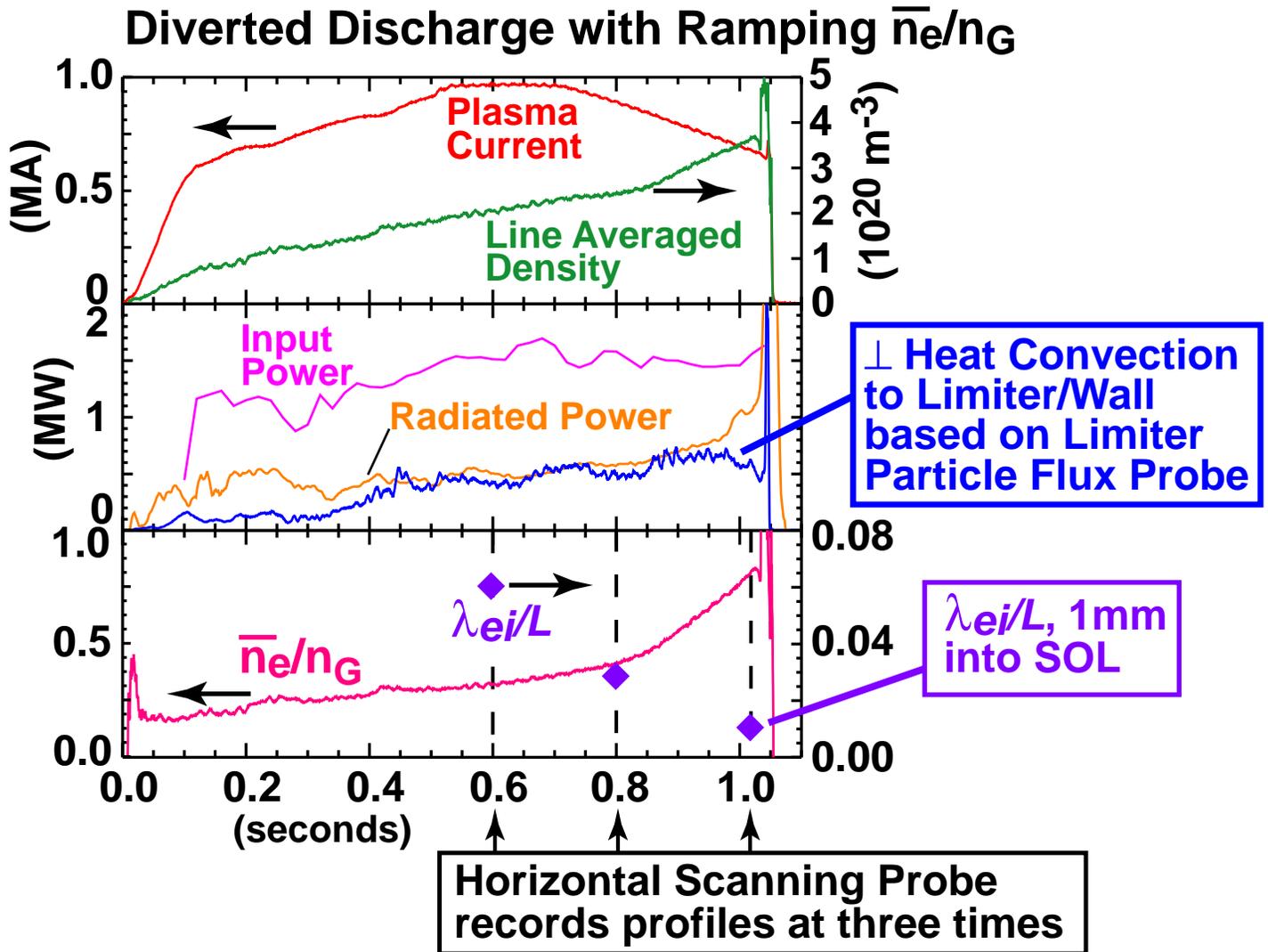
Outline of Talk

- Main-Chamber Particle Balance
- Critical Cross-Field Flux for Particle Control
- Effective Cross-Field Particle Diffusivities (D_{eff}) & Scalings
- Cross-Field Heat Convection
- Character of SOL Fluctuations

- **SOL Transport Physics and the Discharge Density Limit**



Collisionality at the Separatrix and \perp Heat Convection to Limiter/Wall Increases as Discharge Density Limit is Approached



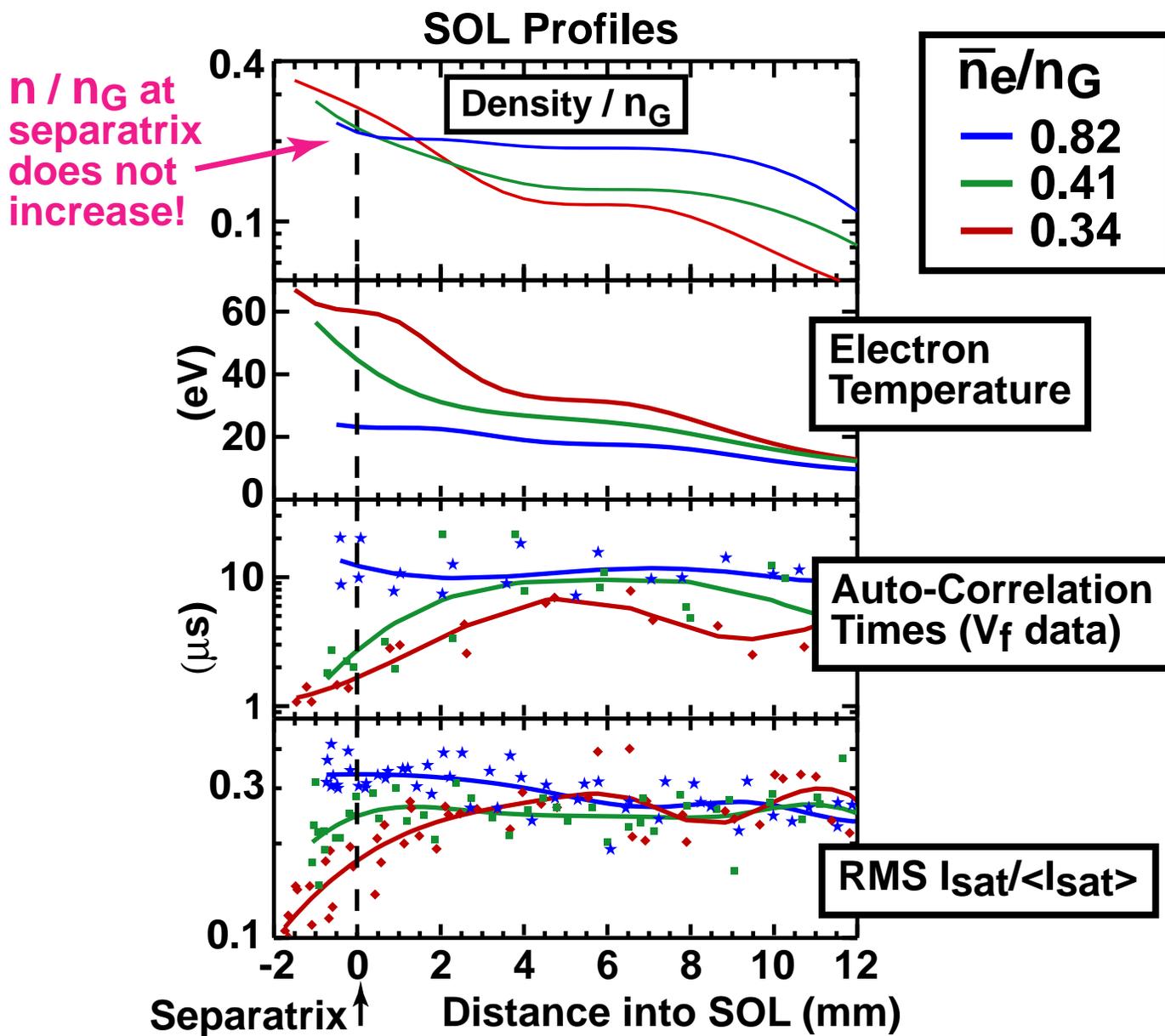
As density limit is approached:

- λ_{ei}/L near separatrix drops dramatically
- Radiation and \perp Convection to Limiter/Wall are comparable and mostly account for input power

Near density limit:

- **Radiation + \perp Convection to Walls \sim Input Power**

Near Density Limit: Large Amplitude, Long-Correlation Time Fluctuations Envelop Entire SOL and Cross Separatrix

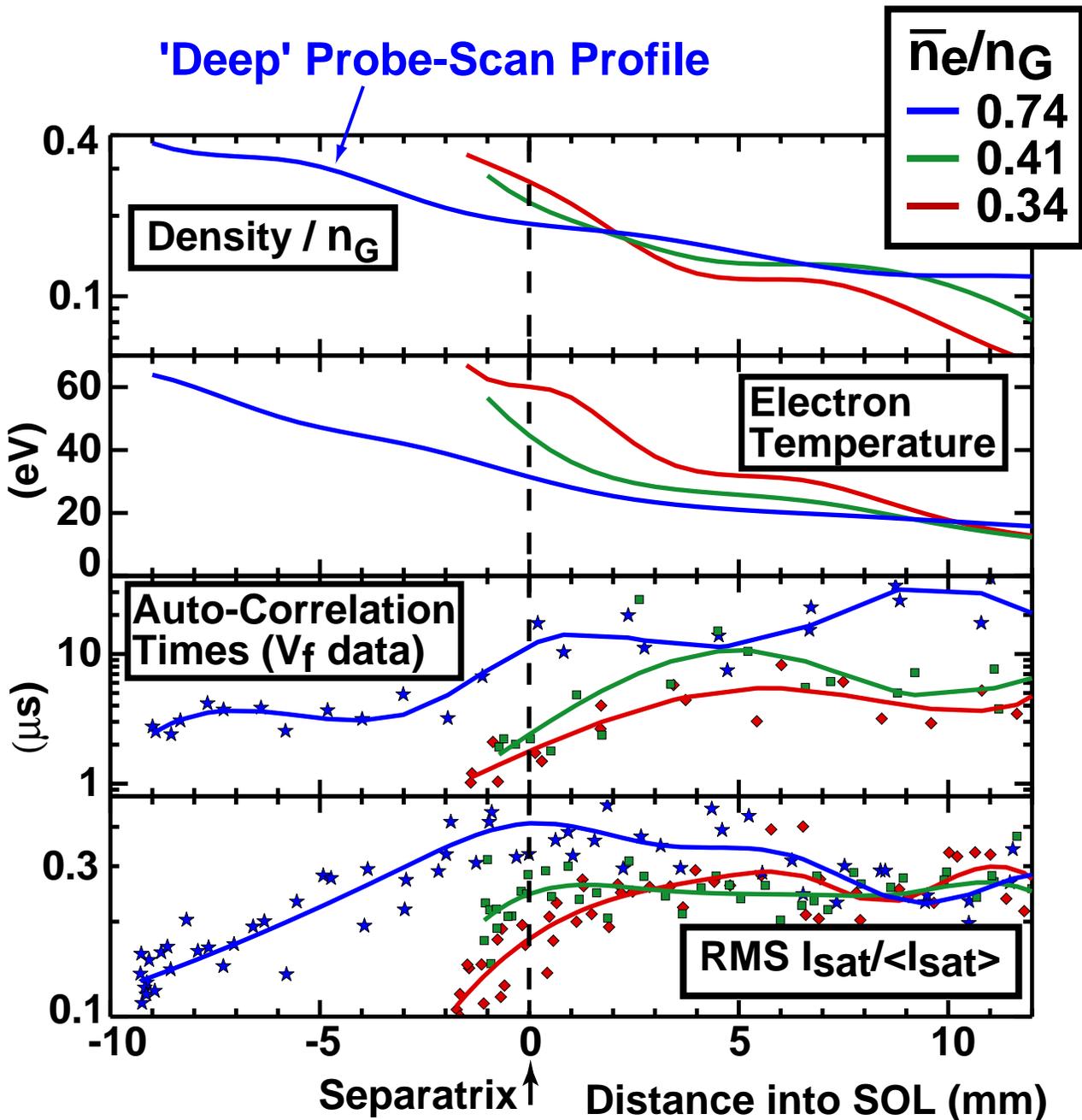


Near density limit:

- SOL n & T_e profiles become flat, T_{sep} low ~ 25 eV!
- Fluctuations characteristic of "Far SOL" now occur everywhere, even across the separatrix

=> Consistent with large \perp Convection Losses

Near Density Limit: 'SOL' Effectively Moves onto Closed Flux Surfaces

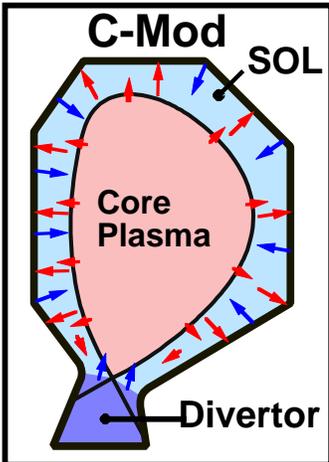


Near density limit:

- Edge n & T_e profiles are apparently *not impacted by magnetic topology!* (open vs closed field lines)

=> Consistent with \perp transport dominating particle and energy losses in edge plasma

Summary



- SOL density profiles exhibit a "two-exponential" decay: *Near and Far SOL*
- Yet, Main-chamber plasma exhausts primarily onto Limiter/Wall surfaces!

- Why? -

=> New Particle Transport Paradigm:

- Density "decays exponentially" because ...
cross-field transport (D_{eff}) increases rapidly with distance into SOL
... not because parallel flows "drain" SOL plasma

- \perp Particle (D_{eff}) and heat convection near separatrix **increases with collisionality:**

$$D_{eff} \sim (\lambda_{ei}/L)^{-1.7}$$

=> Heat Transport Paradigm Modified:

- At moderate collisionality ($\bar{n}_e/n_G \sim 0.5$), cross-field heat convection exceeds conduction losses

T_{sep} no longer regulated by "conduction law":

$$T_{sep} \propto (P_{sol}/\lambda_{Te})^{2/7}$$

Summary (page 2)

- Fluctuation behavior supports picture of particle & energy transport increasing with distance into SOL

Near SOL: (steep density profile) low amplitude "random" fluctuations

Far SOL: (flat density profile) large amplitude intermittent "bursts" in I_{sat} and ~ 1 cm "blobs" in D_{α} , extending into Limiter Shadow

\Rightarrow Large cross-field velocities, > 100 m s $^{-1}$

- $\langle \tilde{n}\tilde{E} \rangle$ -derived cross-field flux (Γ) supports results inferred from particle balance:

Γ similarly increases with \bar{n}_e

Γ is factor of ~ 8 times larger than particle balance
 \Rightarrow supporting MCR conclusion

- Far SOL turbulence has some SOC-like characteristics (\sim generic to edge plasmas)

PDF(Γ) has power-law tail (independent of \bar{n}_e)

Summary (page 3)

=> New Insight on Density Limit Physics:

- As density limit is approached, λ_{ei}/L near separatrix drops and transport across the SOL increases dramatically

⊥ Heat Convection to Limiter/Wall becomes large fraction of input power

"Bursty" fluctuations (large transport) occur over entire SOL and begins to attack plasma on closed flux surfaces

~ at limit:

Radiation + Convection to Wall ~ Input Power

Rapid increase of ⊥ Heat Convection as edge plasma cools may play role in thermal instability leading to disruption



Need to develop scaleable empirical and physics-based understandings of underlying transport physics